



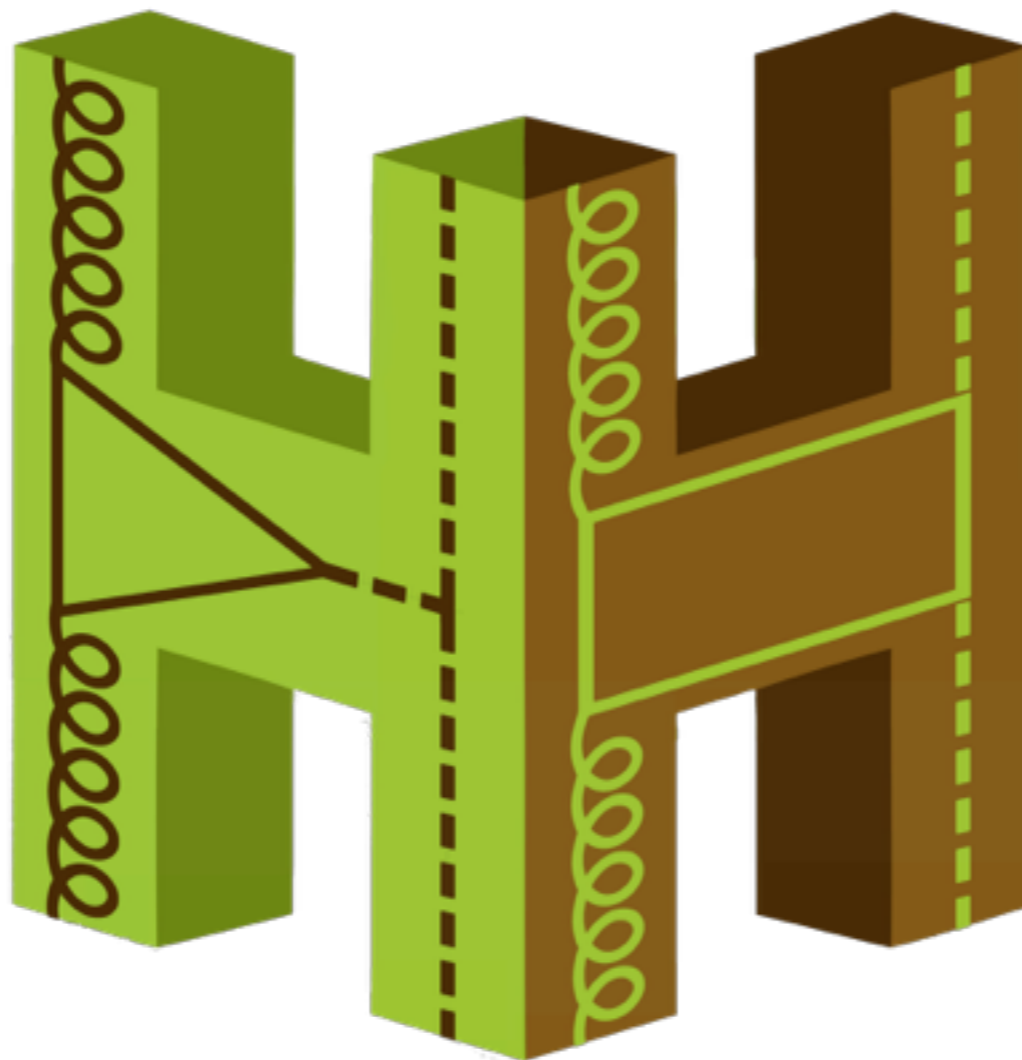
Searching for Di-Higgs \rightarrow $b\bar{b}\tau\tau$ at ATLAS

Katharine Leney

20th November 2015

Overview

- Motivation for searching for di-Higgs pairs (and in the $bb\tau\tau$ channel).
- ATLAS Run 1 $HH \rightarrow bb\tau\tau$ analysis.
- Combination of all ATLAS HH channels.
- Prospects for $HH \rightarrow bb\tau\tau$ searches in Run 2 and beyond.

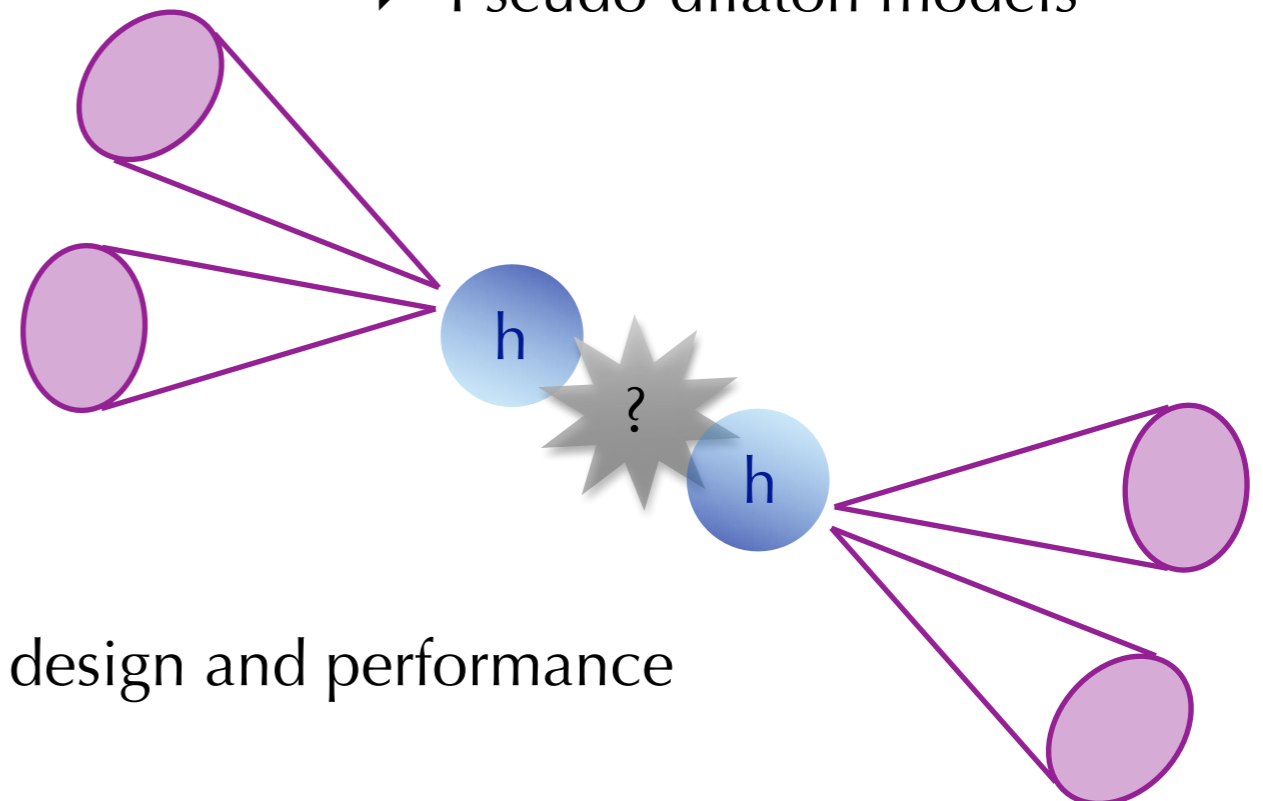


Motivation

- TeV-scale resonances decaying to two 125 GeV Higgs bosons (h) predicted by several models, including:
 - ▶ RS KK Graviton
 - ▶ 2 Higgs Doublet Models (2HDM)
 - ▶ Higgs portal models
 - ▶ Composite models with hh resonances



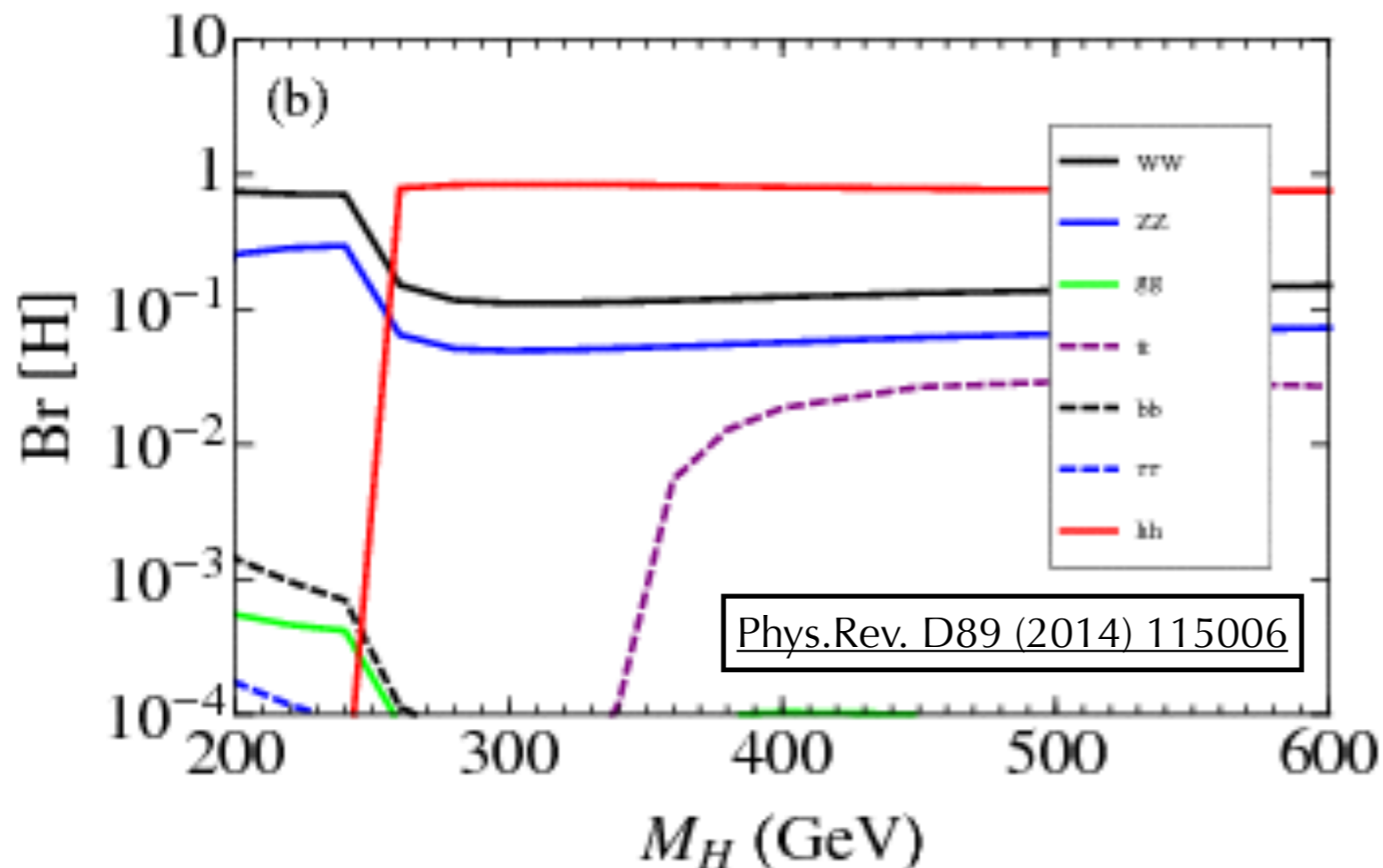
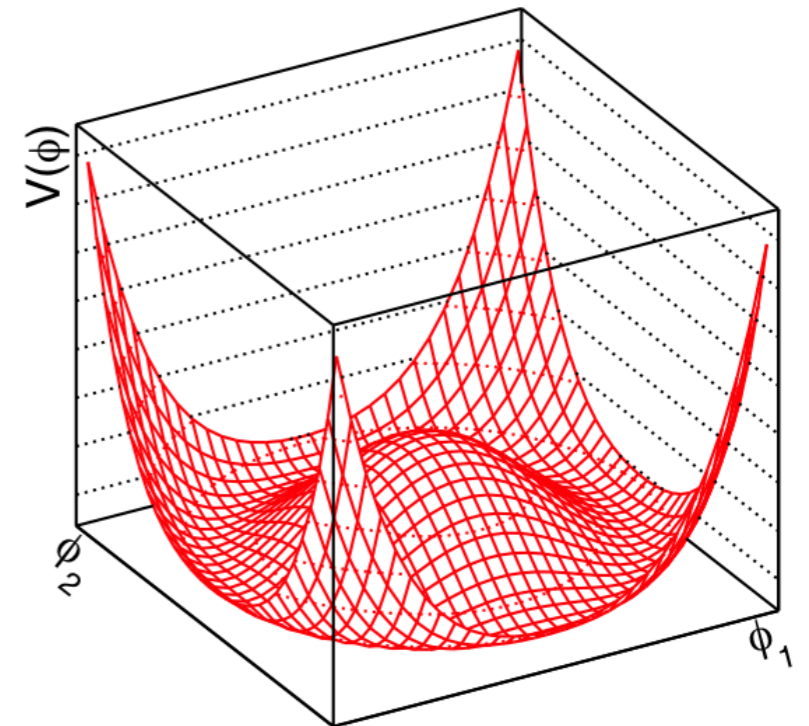
- Enhancement of non-resonant di-Higgs production, e.g.
 - ▶ Models with heavy top-partners
 - ▶ Composite Higgs models
 - ▶ Pseudo-dilaton models



- SM di-Higgs production at HL-LHC.
 - ▶ Need $\sim 3000 \text{ fb}^{-1}$ to measure this.
 - ▶ Sensitivity studies now drive upgrade design and performance requirements.

SUSY & The Higgs Sector (2HDM)

- Supersymmetric extensions to Standard Model require two Higgs doublets \rightarrow 5 observable Higgs bosons:
 - 3 neutral (h/A/H)
 - 2 charged (H^\pm)
- Assume the observed 125 GeV Higgs boson is the light Higgs (h).
- Heavy Higgs (H) can decay to a pair of light Higgses (h).



- Branching ratio to hh can be large, depending on parameters of model.
- Range $m_H < 500$ GeV theoretically favoured.
- Couplings can be expressed as functions of α (mixing angle) and $\tan \beta$ (ratio of vacuum expectation values).

Di-Higgs Decay Channels

	bb	WW	$\tau\tau$	ZZ	$\gamma\gamma$
bb	33%				
WW	25%	4.6%			
$\tau\tau$	7.4%	2.5%	0.39%		
ZZ	3.1%	1.2%	0.34%	0.076%	
$\gamma\gamma$	0.26%	0.10%	0.029%	0.013%	0.0053%

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bbbb

- Largest branching ratio
- Harder to trigger on (especially for low mass resonances and SM)

Di-Higgs Decay Channels

	bb	WW	$\tau\tau$	ZZ	$\gamma\gamma$
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xxWW,

- Take a hit on branching ratio to have events where W decays leptonically (~34%).
- Reconstructing hadronic W decays difficult (large QCD backgrounds).

Di-Higgs Decay Channels

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xxYY

- Clean signature (thanks to di-photon pair in final state).
- Tiny branching ratios...

Di-Higgs Decay Channels

	bb	WW	$\tau\tau$	ZZ	$\gamma\gamma$
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ZZZZ

- Possibility to select clean final states with 2 or 4 leptons, but large hit on branching ratio.

Di-Higgs Decay Channels

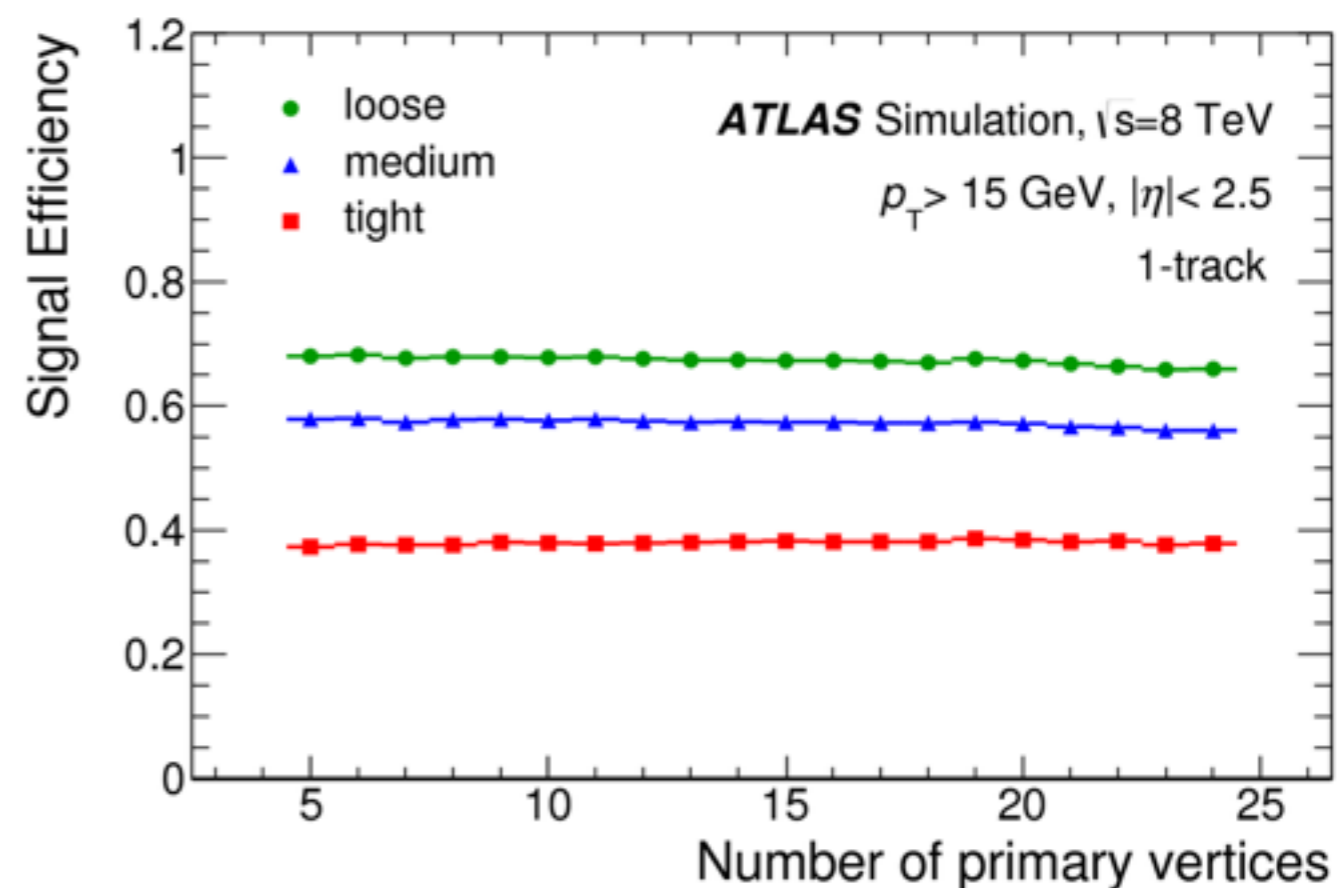
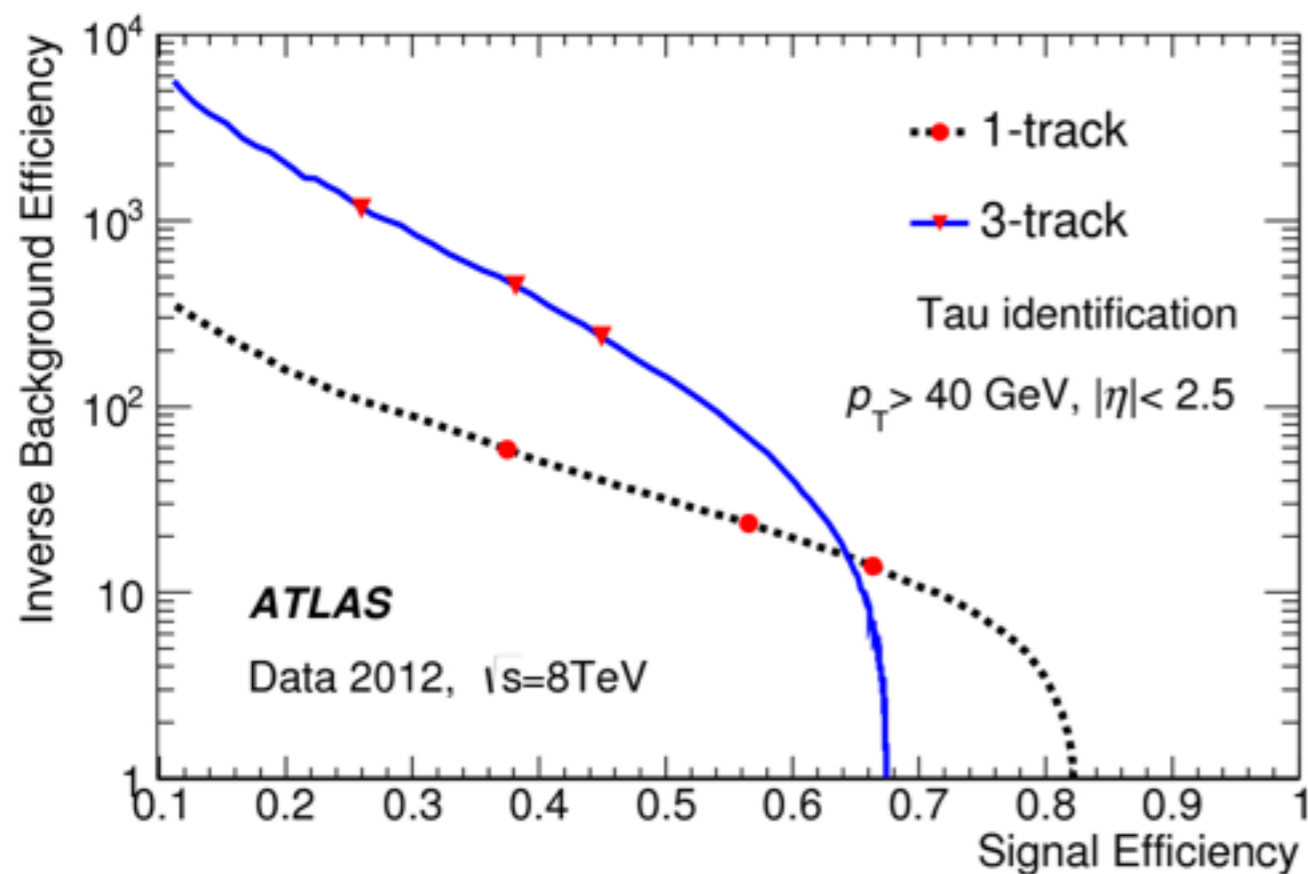
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bb $\tau\tau$

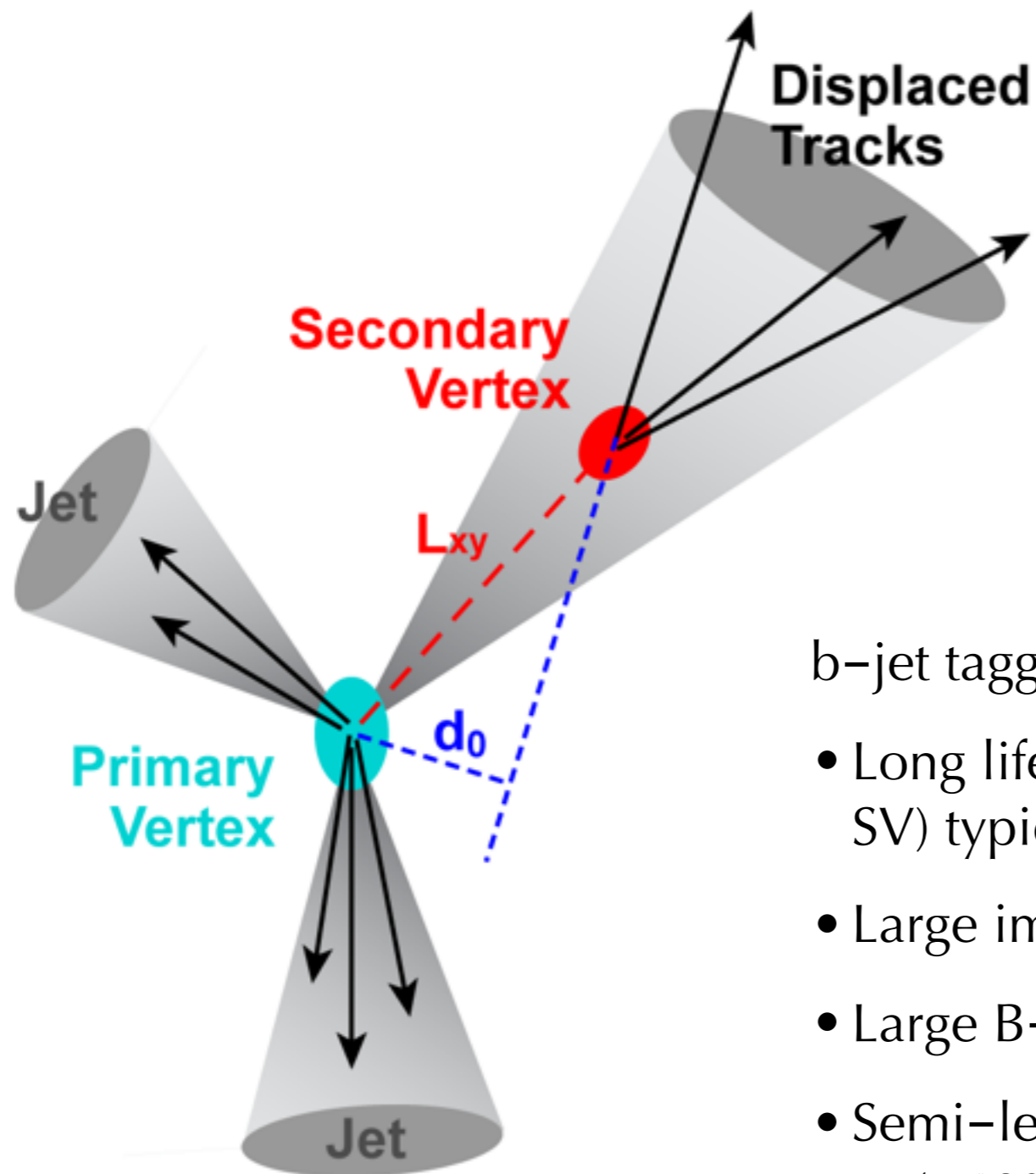
- Possibility to trigger on lepton from tau decays (58% of events).
- Fully hadronic tau-tau channels have comparable sensitivity to $e\tau_h$ and $\mu\tau_h$ channels at ATLAS.

Tau Reconstruction and ID in ATLAS

- Taus seeded from anti- k_t jets with $\Delta R = 0.4$.
- Candidates required to be associated with 1 or 3 tracks within a core region $\Delta R < 0.2$.
- Isolated in annulus $0.2 < \Delta R < 0.4$.
- Various discriminating variables combined in Boosted Decision Trees to reject tau-fakes from electrons and jets.
 - ▶ BDTs trained separately for 1 and 3-prong taus.



b-Jet Decay Characteristics

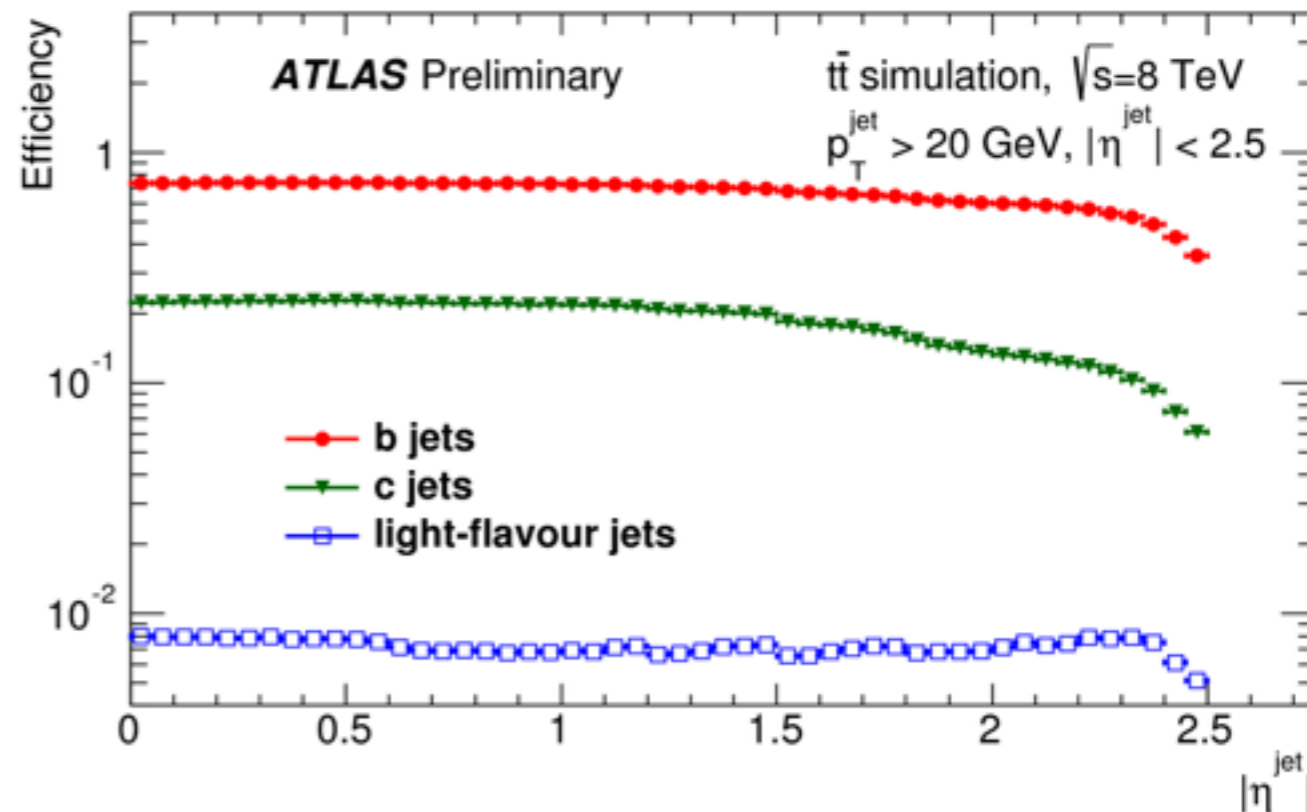
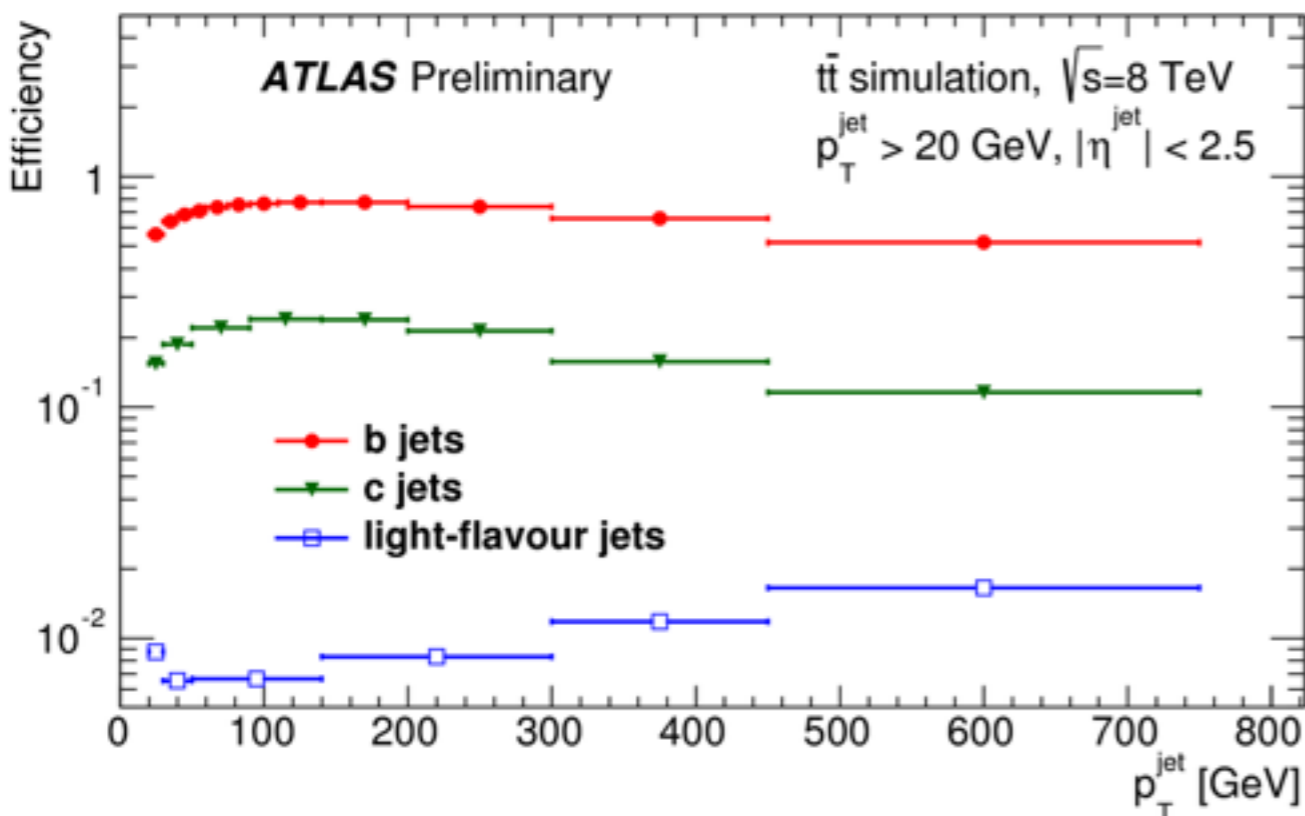
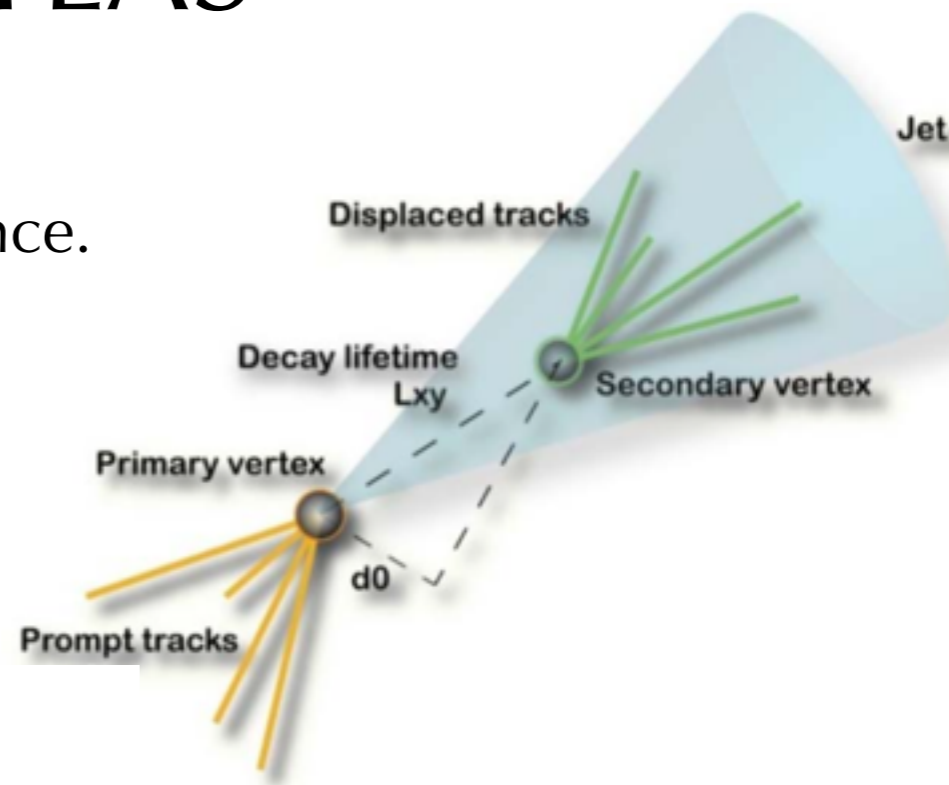


b-jet tagging relies on B-hadron properties:

- Long lifetime \rightarrow displaced vertex (secondary vertex, SV) typically few mm from primary vertex (PV).
- Large impact parameter (d_0).
- Large B-hadron mass.
- Semi-leptonic (e/μ) decay of B-hadron
 - ($\sim 40\%$ including $b \rightarrow c \rightarrow \ell \nu X$ decay).

b-Tagging in ATLAS

- Use outputs of 3 algorithms as inputs to MVA:
 - ▶ IP3D: Use transverse and longitudinal IP significance.
 - ▶ SV1: Reconstruct SV and use information about:
 - SV mass
 - $\Sigma(P_T \text{ SV tracks}) / \Sigma(P_T \text{ all tracks in jet})$
 - Number of two-track vertices
 - ΔR (jet-direction, PV \rightarrow SV direction)
 - ▶ JetFitter: Exploit the topology of weak B/C-hadron decay chain ($b \rightarrow c \rightarrow X$) inside jets.



Run 1 HH \rightarrow bb $\tau\tau$ Analysis

Preselection

- Select all objects in bb $\tau\tau$ final state ($\ell + \tau + 2$ -jets).
- Only lepton-hadron decay mode considered in Run 1.

Phys. Rev. D 92,
092004 (2015)

Kinematic Cuts

- Addition rejection against Z+jets, W+jets and ttbar.

Event Categorisation

- Number of b-jets.
- p_T of τ -pair.

Resonant Search

- $m(\tau\tau)$ cut
- Final discriminant: $m(bb\tau\tau)$.

Non-Resonant Search

- Final discriminant: $m(\tau\tau)$

$\sqrt{s} = 8 \text{ TeV}$
 $\int \mathcal{L} = 20 \text{ fb}^{-1}$

Event Preselection

Single lepton trigger

One isolated lepton (e/ μ), $p_T > 26 \text{ GeV}$

Di-lepton veto

One hadronic tau, $p_T > 20 \text{ GeV}$

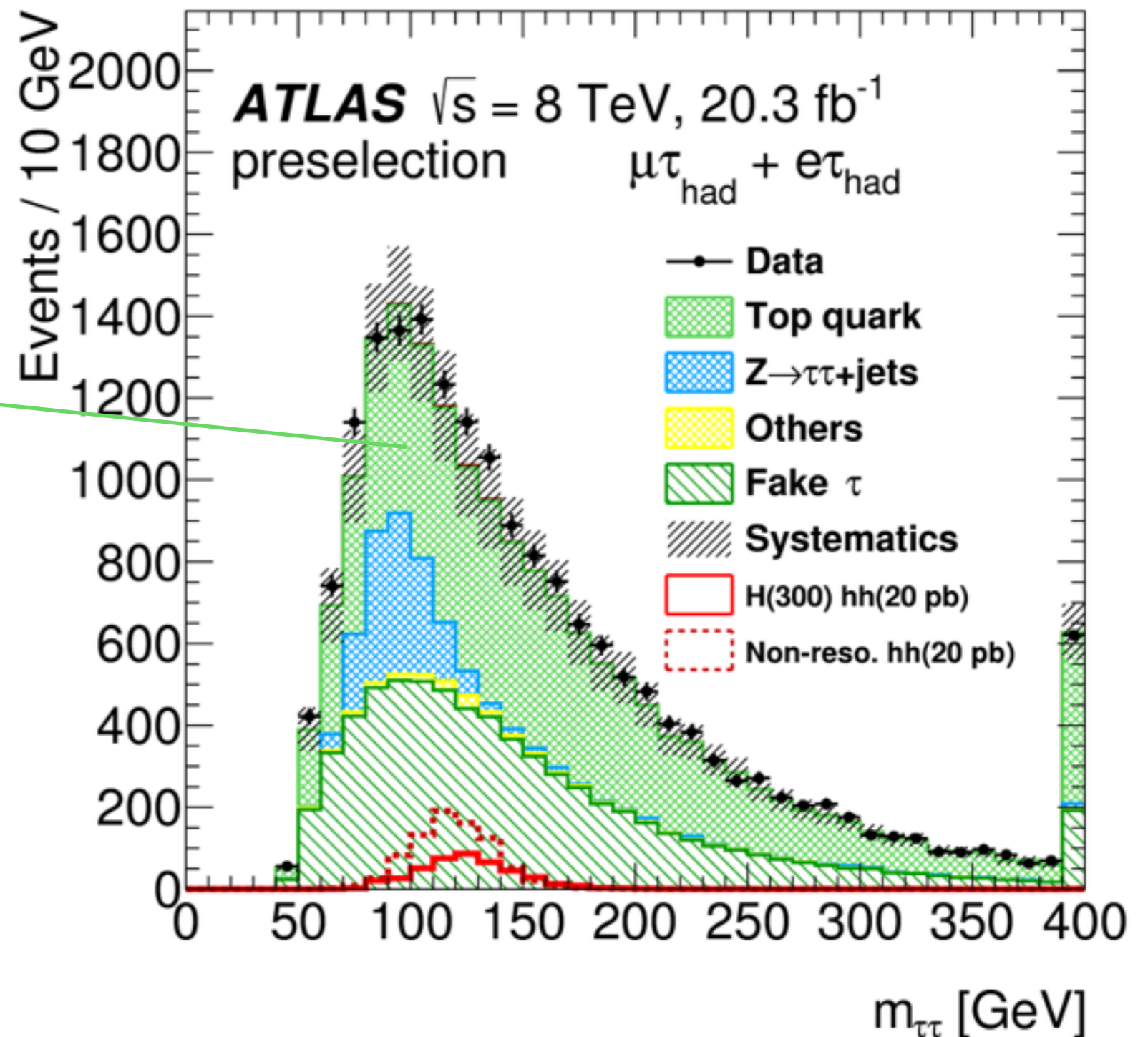
Charge correlation between lepton and tau

Two or more jets, $p_T > 30 \text{ GeV}$

Backgrounds

Top-Pair Production

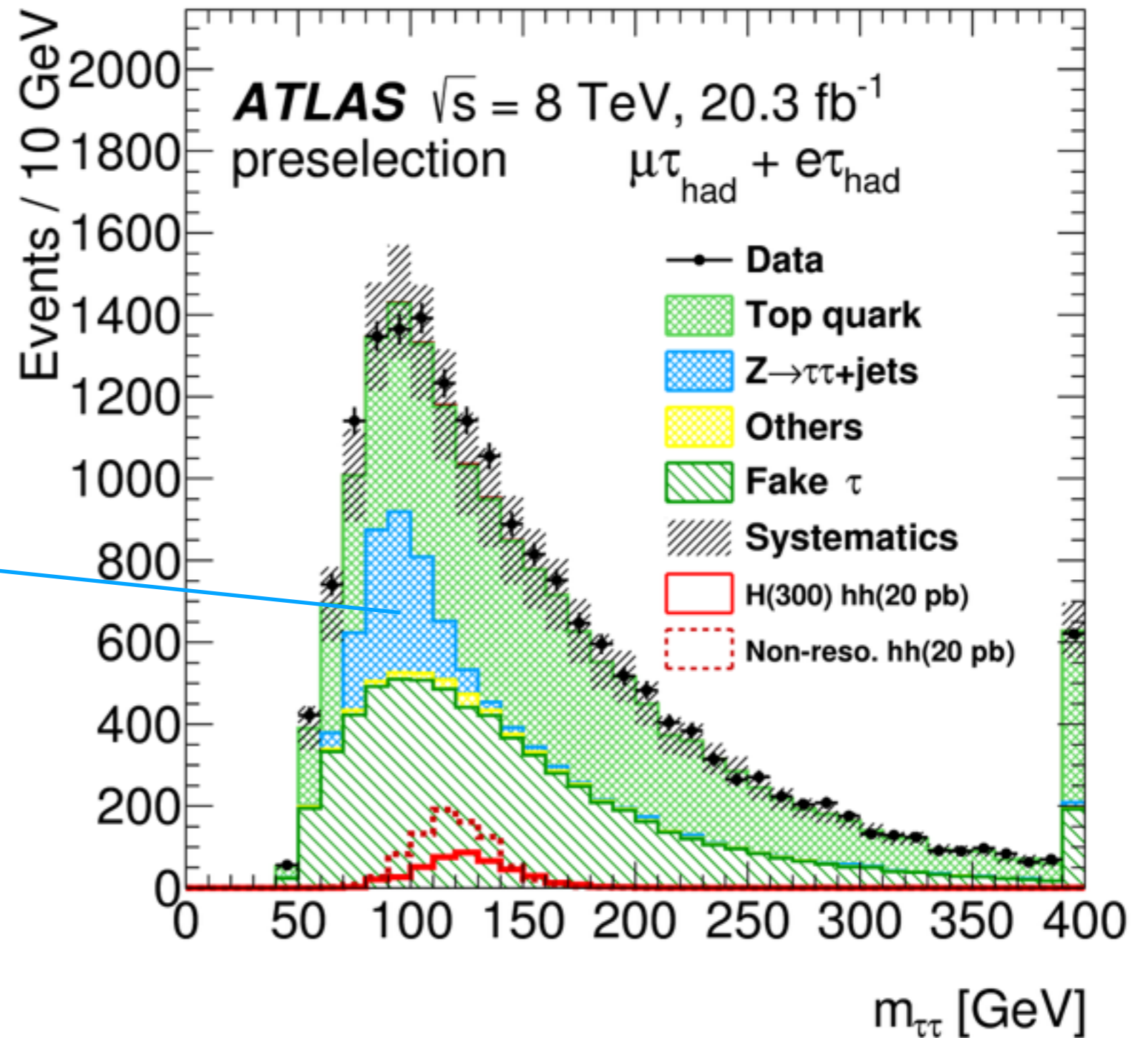
- Contribution where the hadronic tau is real estimated from MC (Powheg).
- Fraction where hadronic W-decay fakes tau calculated separately (see later slide).



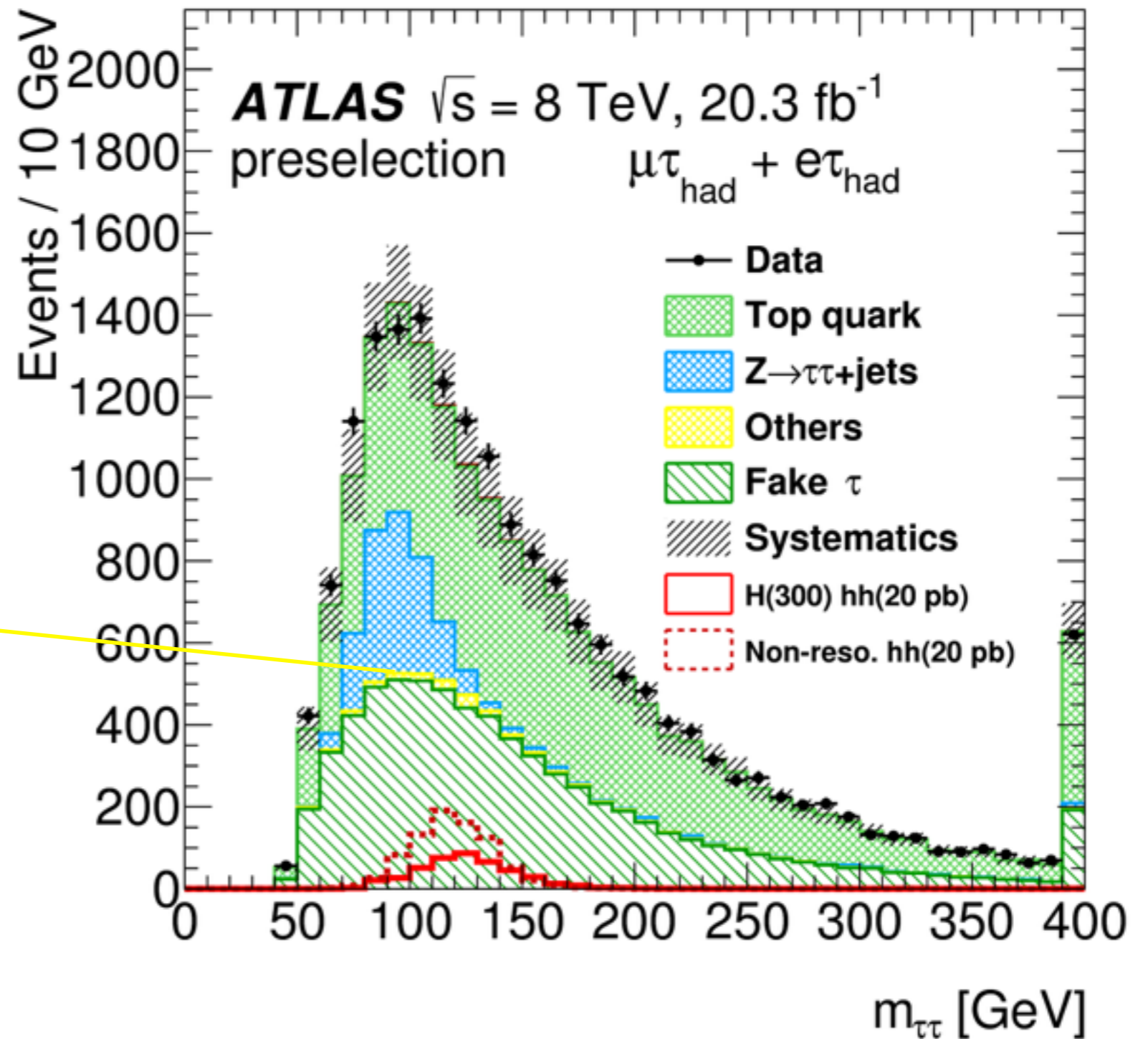
Backgrounds

$Z \rightarrow \tau\tau$ Background

- Estimated from data using embedding method:
 - ▶ $Z \rightarrow \mu\mu$ events selected from data.
 - ▶ Muons replaced with simulated taus.
 - ▶ Missing E_T corrections applied.
- Normalised to data in $40 < M^{\text{vis}}(\tau\tau) < 70$ GeV region.



Backgrounds



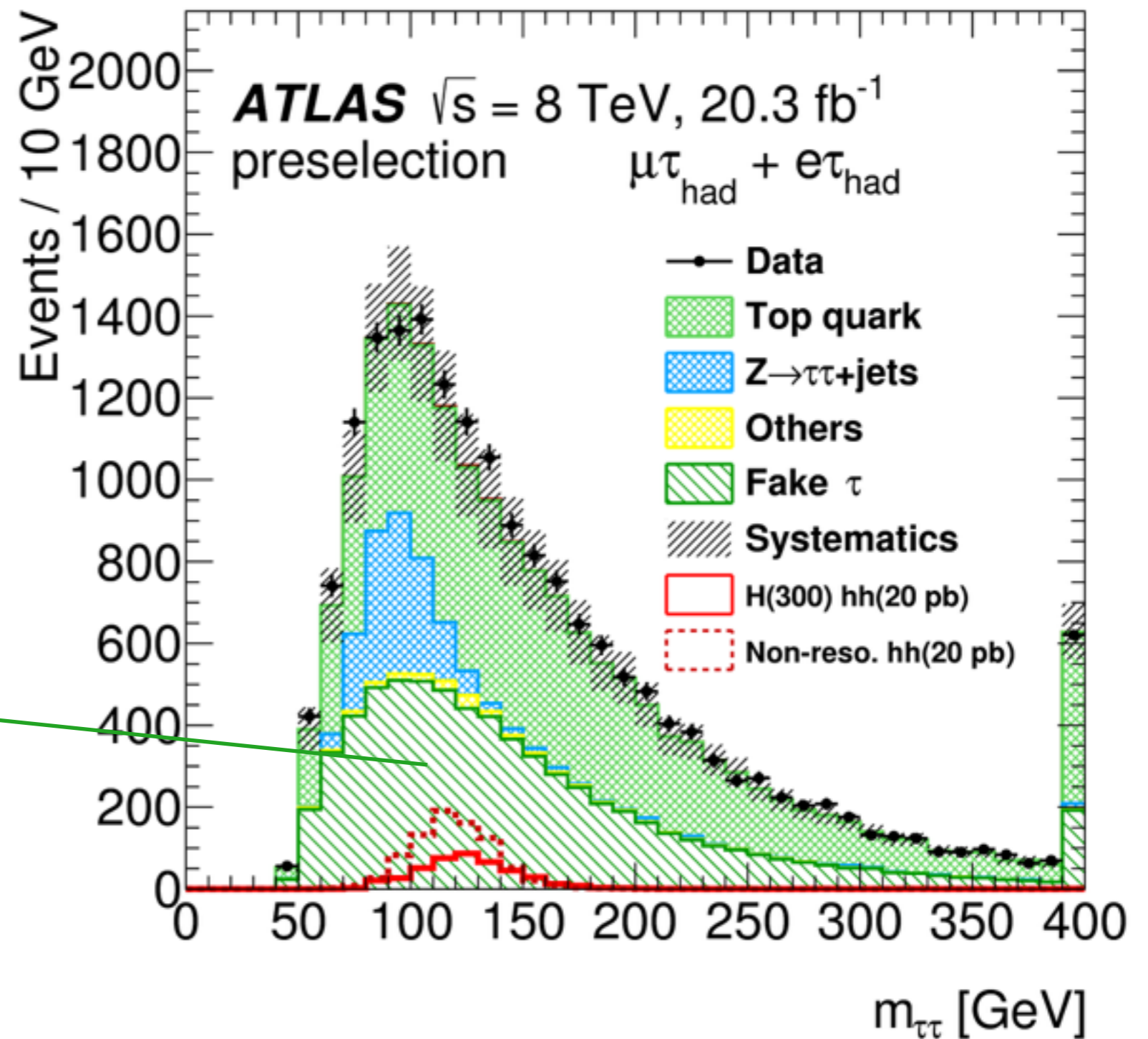
Others

- Single top, di-boson, $Z \rightarrow \ell\ell$ all estimated from MC.
- Small contribution.

Backgrounds

Fake Taus

- Fake-factor method to estimate all backgrounds where the tau is faked by a jet.
 - ▶ Multijets.
 - ▶ Semi-leptonic $t\bar{t}$.
 - ▶ W +jets.
 - ▶ $Z \rightarrow \tau\tau$ +jets (where one tau is missed).



Fake-Factor Method

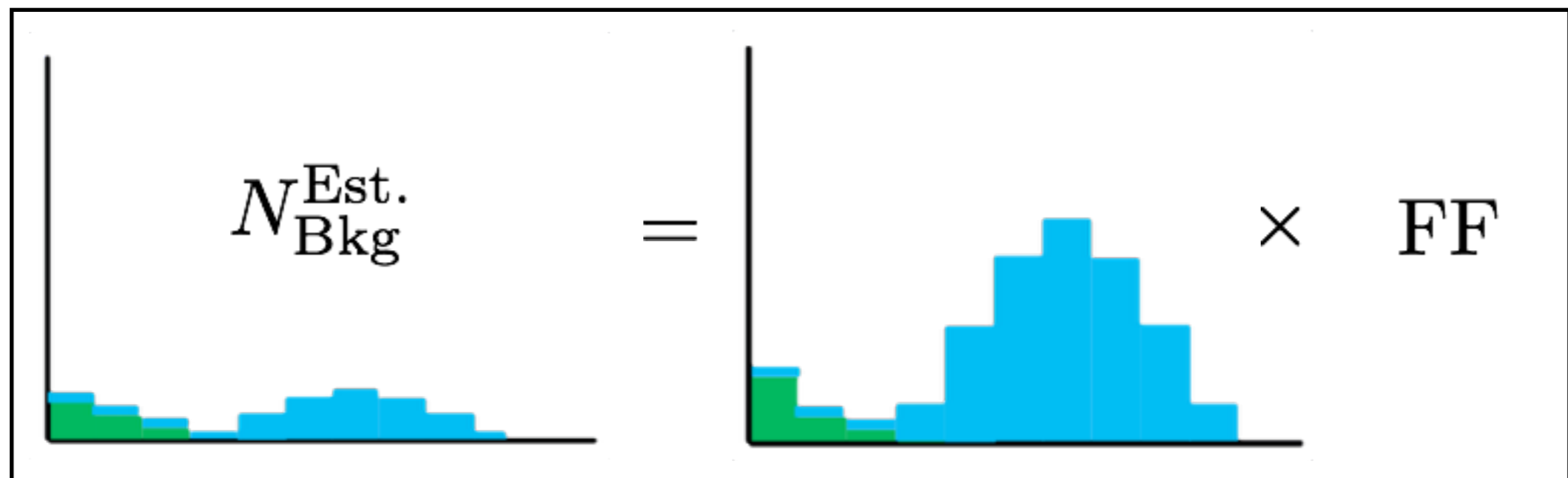
$$N_{\text{Bkg}}^{\text{Est.}} = (N_{\text{data,SR}}^{\text{anti-}\tau} - N_{\text{true,SR}}^{\text{anti-}\tau}) \times \text{FF}_{\text{CR}} \quad \text{FF}_{\text{CR}} = \frac{N_{\text{CR}}^{\text{identified-}\tau}}{N_{\text{CR}}^{\text{anti-}\tau}}$$

- Number of anti-tau events (failing signal tau ID, but passing looser selection).
- Contribution from real taus subtracted.

- Separate FFs for each contributing process.
- Calculated as a function of p_{T} and #-tracks.

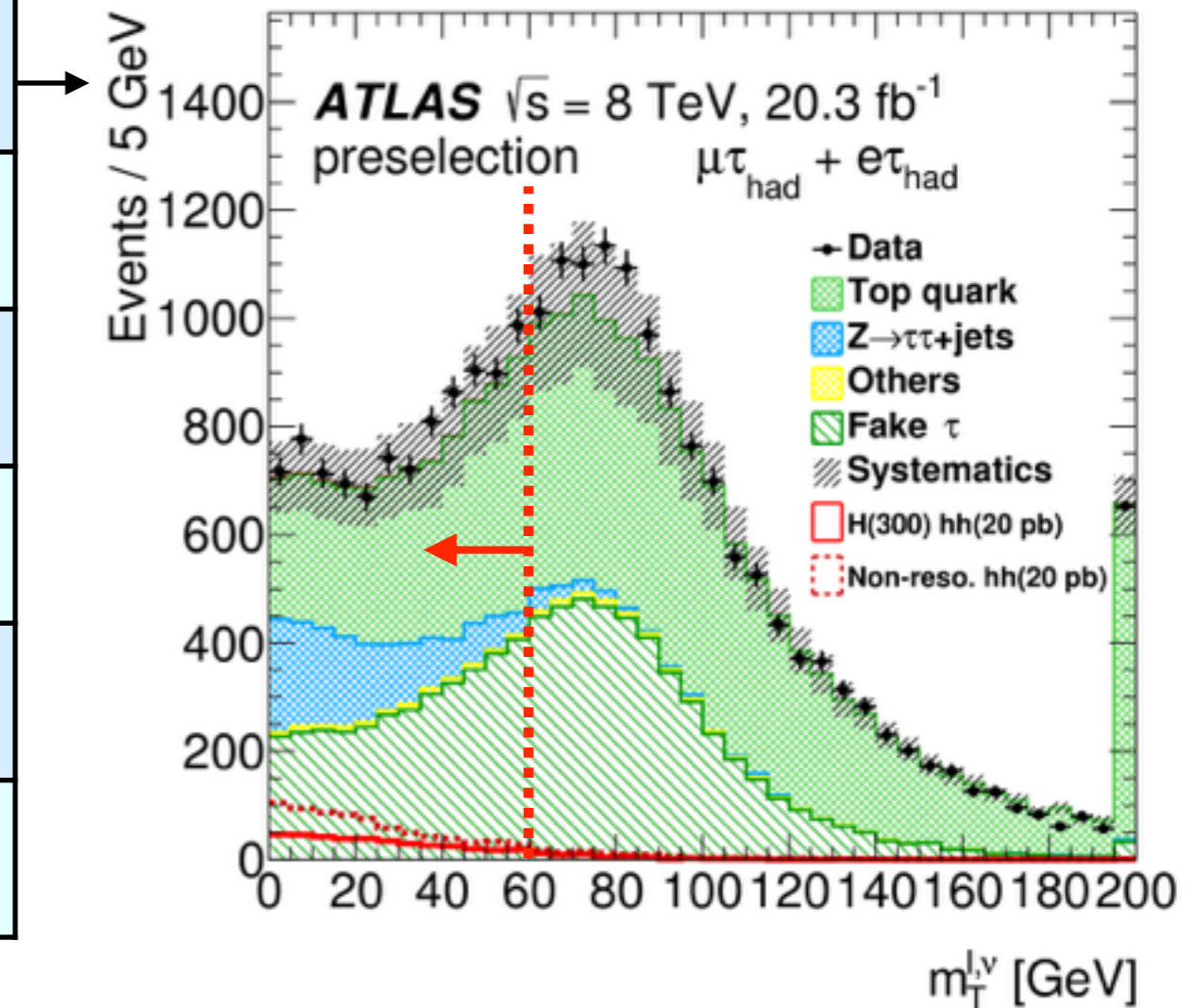
$$\text{FF}(p_{\text{T}}, n_{\text{prong}}) = \sum_{i=\text{bkg}} R_i \text{FF}_i(p_{\text{T}}, n_{\text{prong}})$$

R_i = fraction of each source
 FF_i = FF of each source



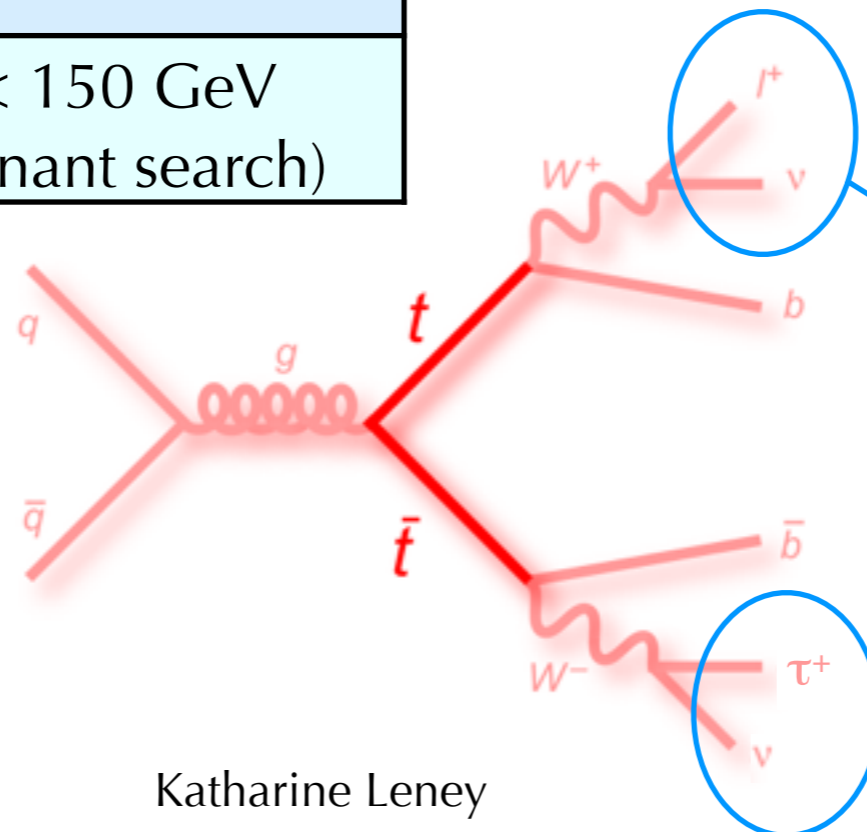
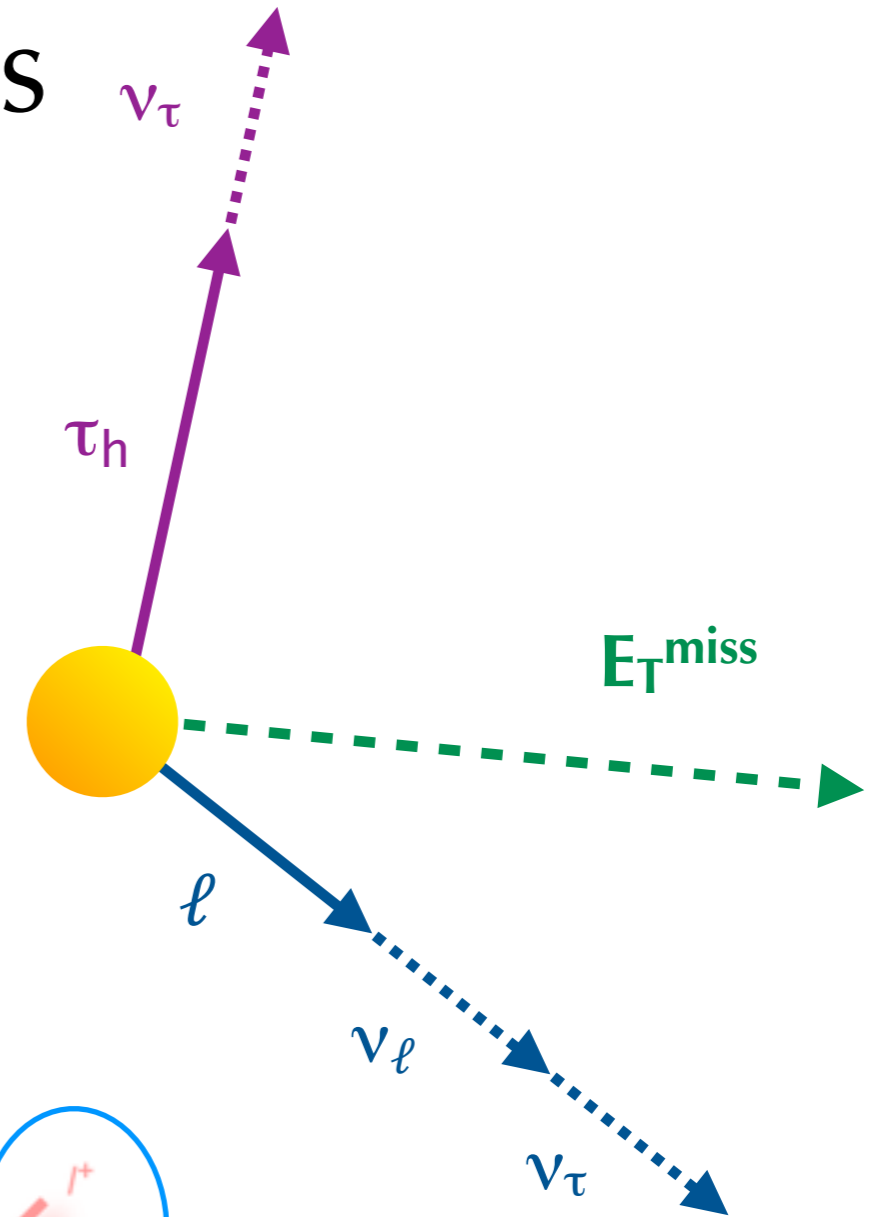
Kinematic Cuts

Transverse mass	Reject W +jets and $t\bar{t}$ $m_T < 60$ GeV
Missing E_T ϕ centrality	Missing E_T should point between the lepton and tau.
Lepton-tau p_T balance	$\Delta p_T(\ell, \tau) < 20$ GeV
m_W vs. m_{top}	Elliptical cuts on W and top masses
m_{bb} mass window	$90 < m_{bb} < 160$ GeV
$m_{\tau\tau}$ mass window	$100 < m_{\tau\tau} < 150$ GeV (only for resonant search)



Kinematic Cuts

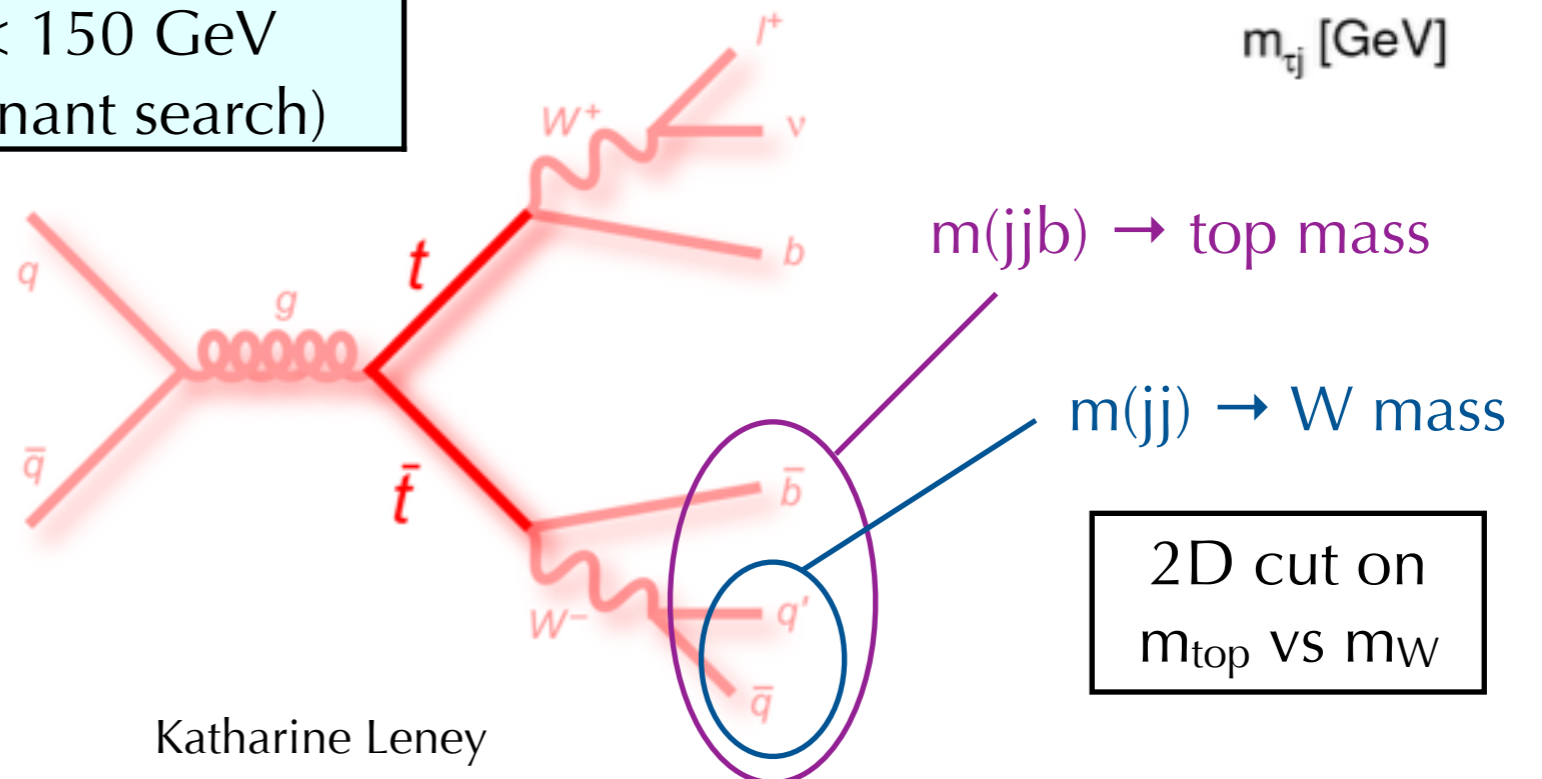
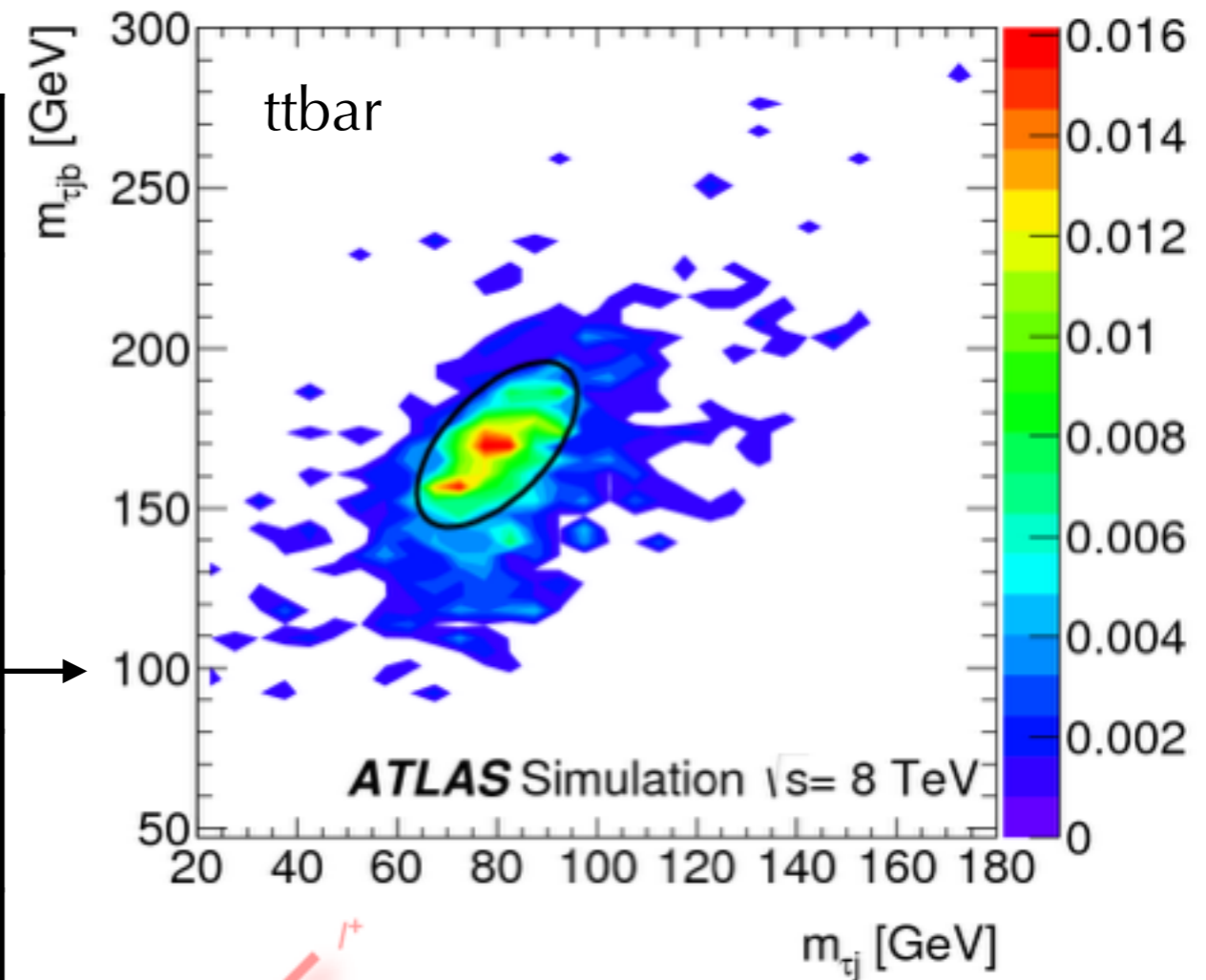
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Leptons from W decay tend to be harder than those from taus because of number of accompanying neutrinos

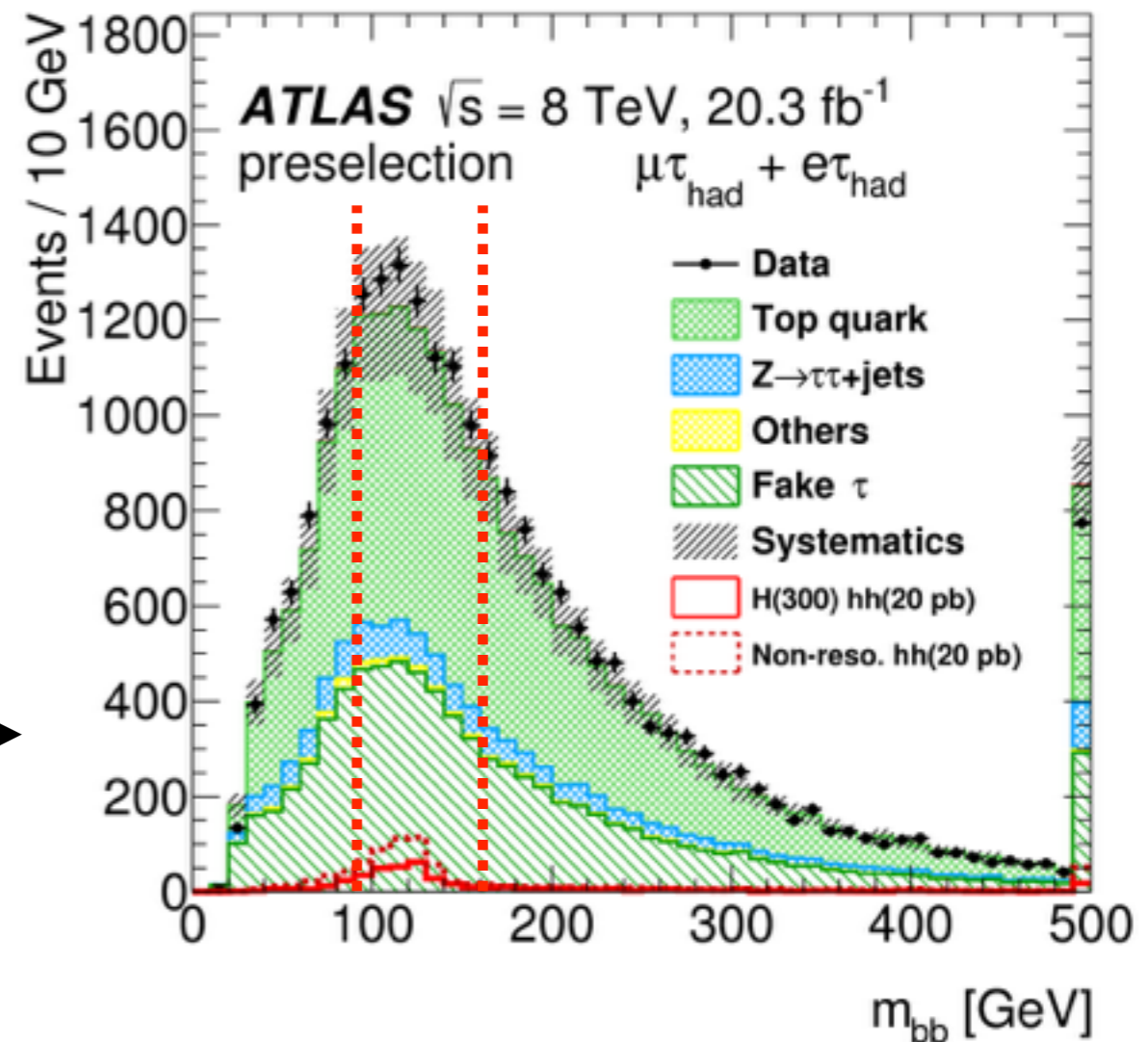
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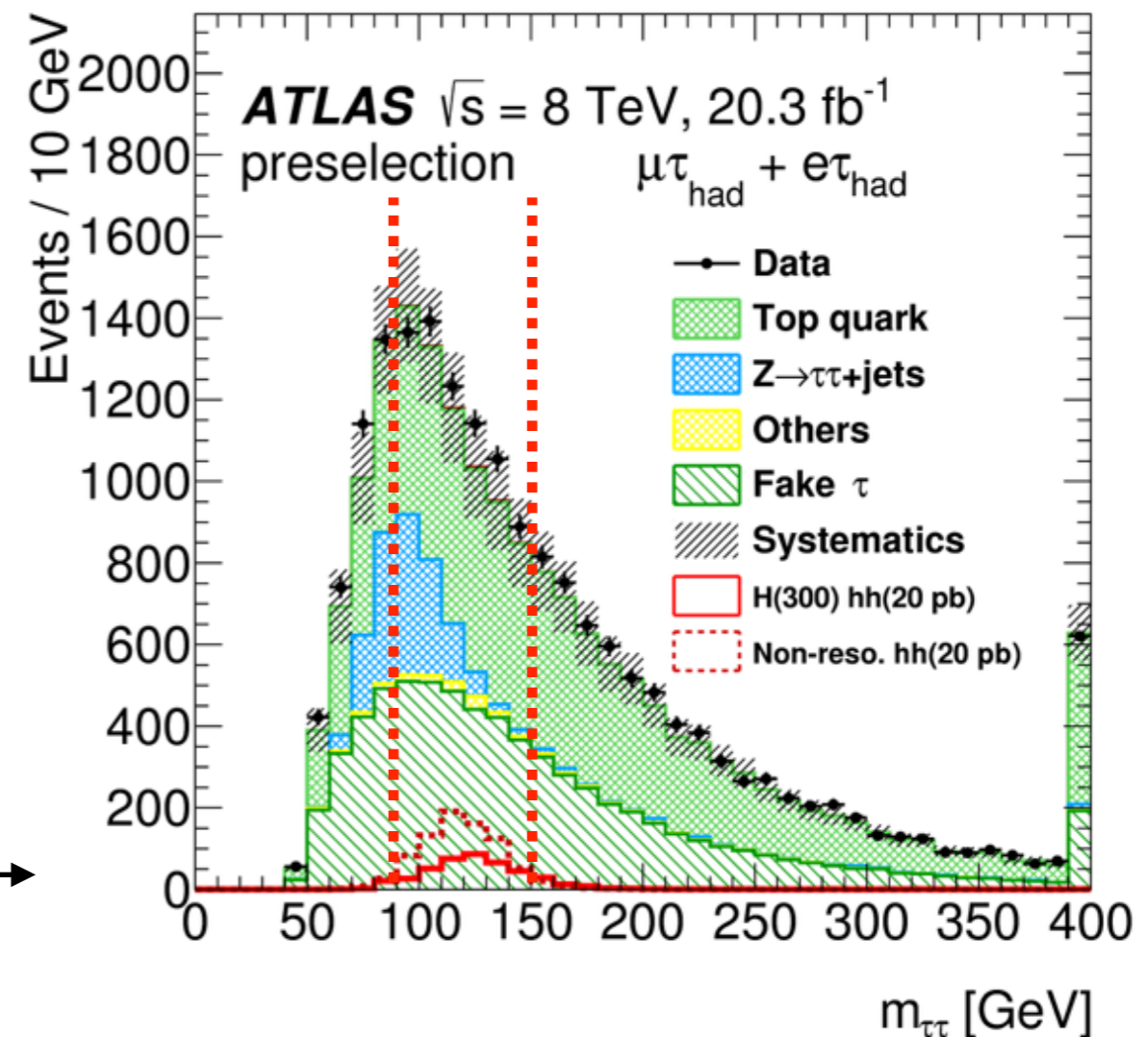
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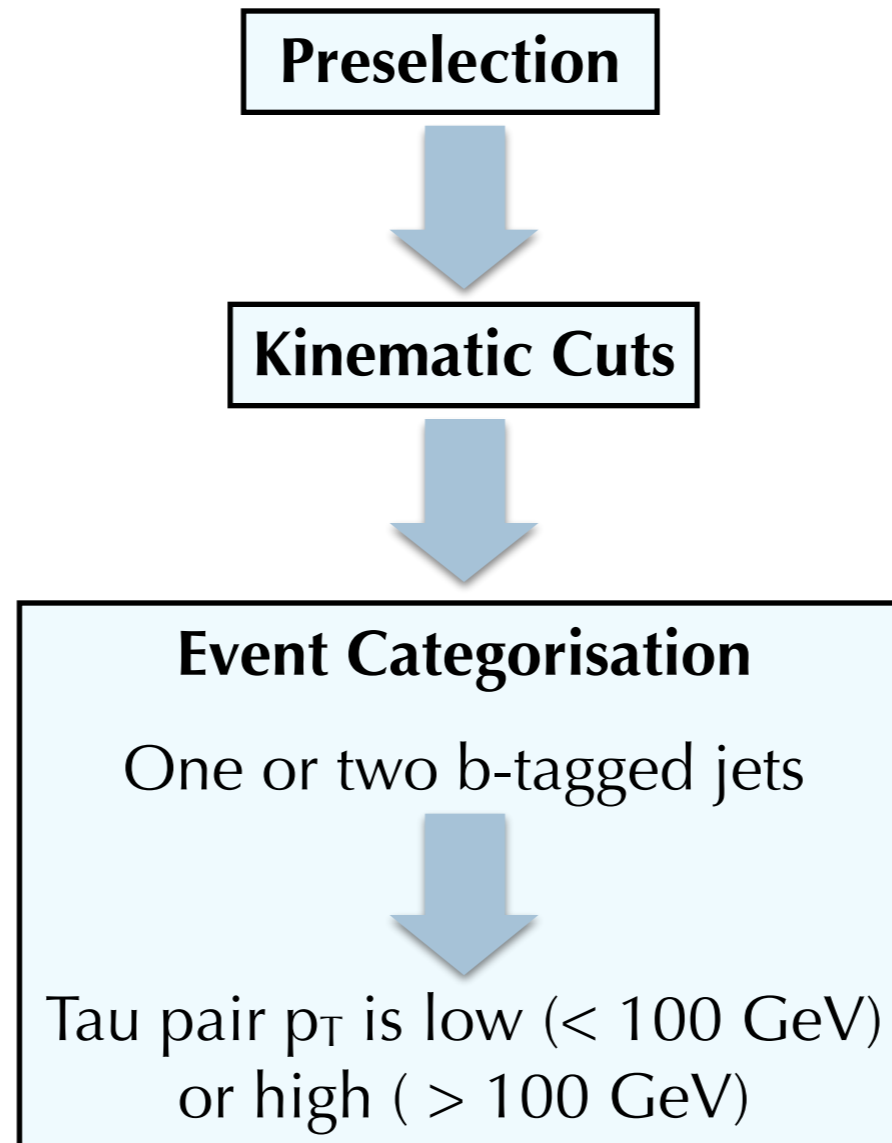
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- $m_{\tau\tau}$ reconstructed using Missing Mass Calculator (MMC), arXiv: 1012.4686
- Weights the kinematically allowed τ decay solutions by a likelihood function.

Event Categorisation

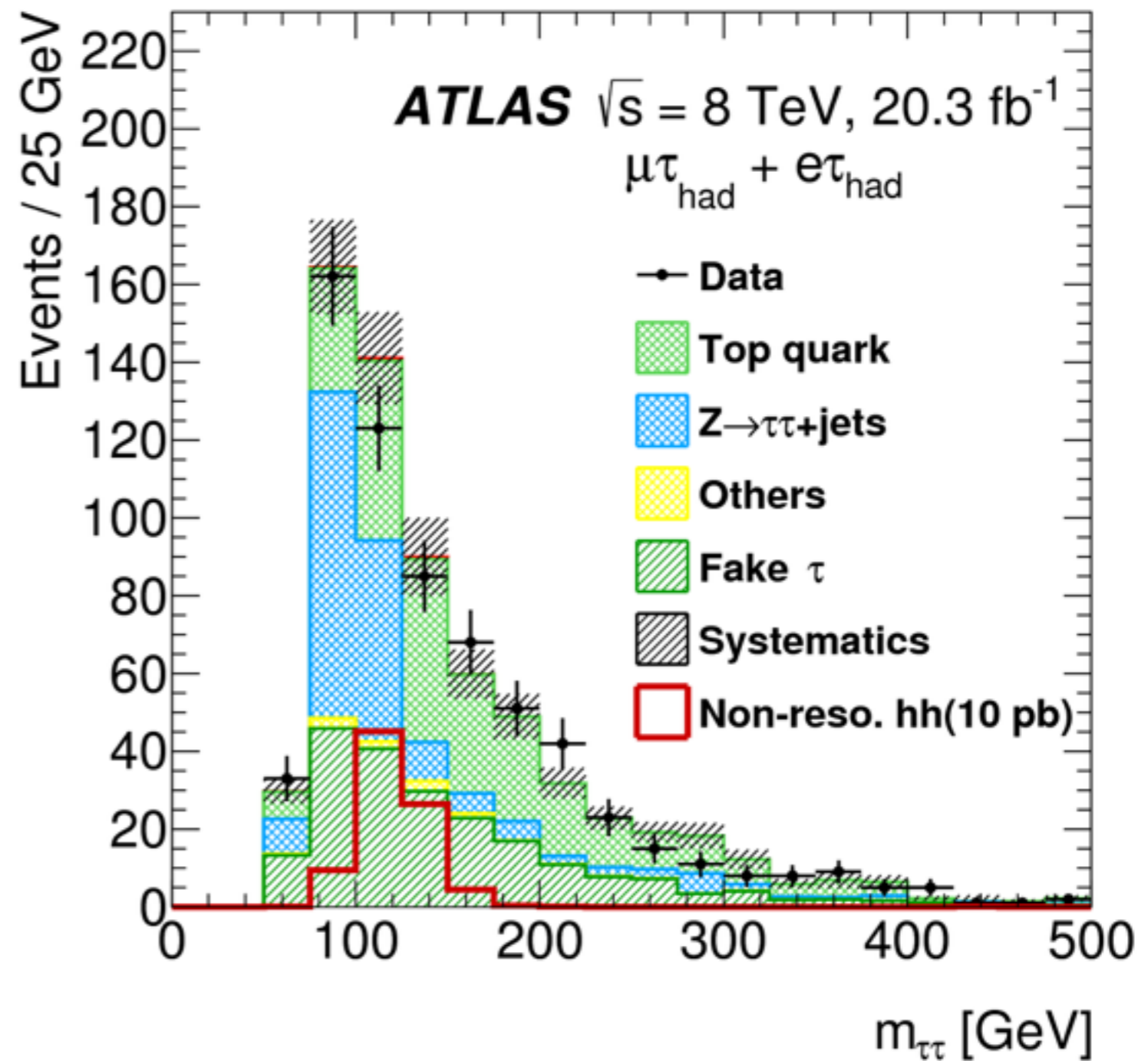


Non-Resonant Analysis

- Use $m_{\tau\tau}$ as final discriminant

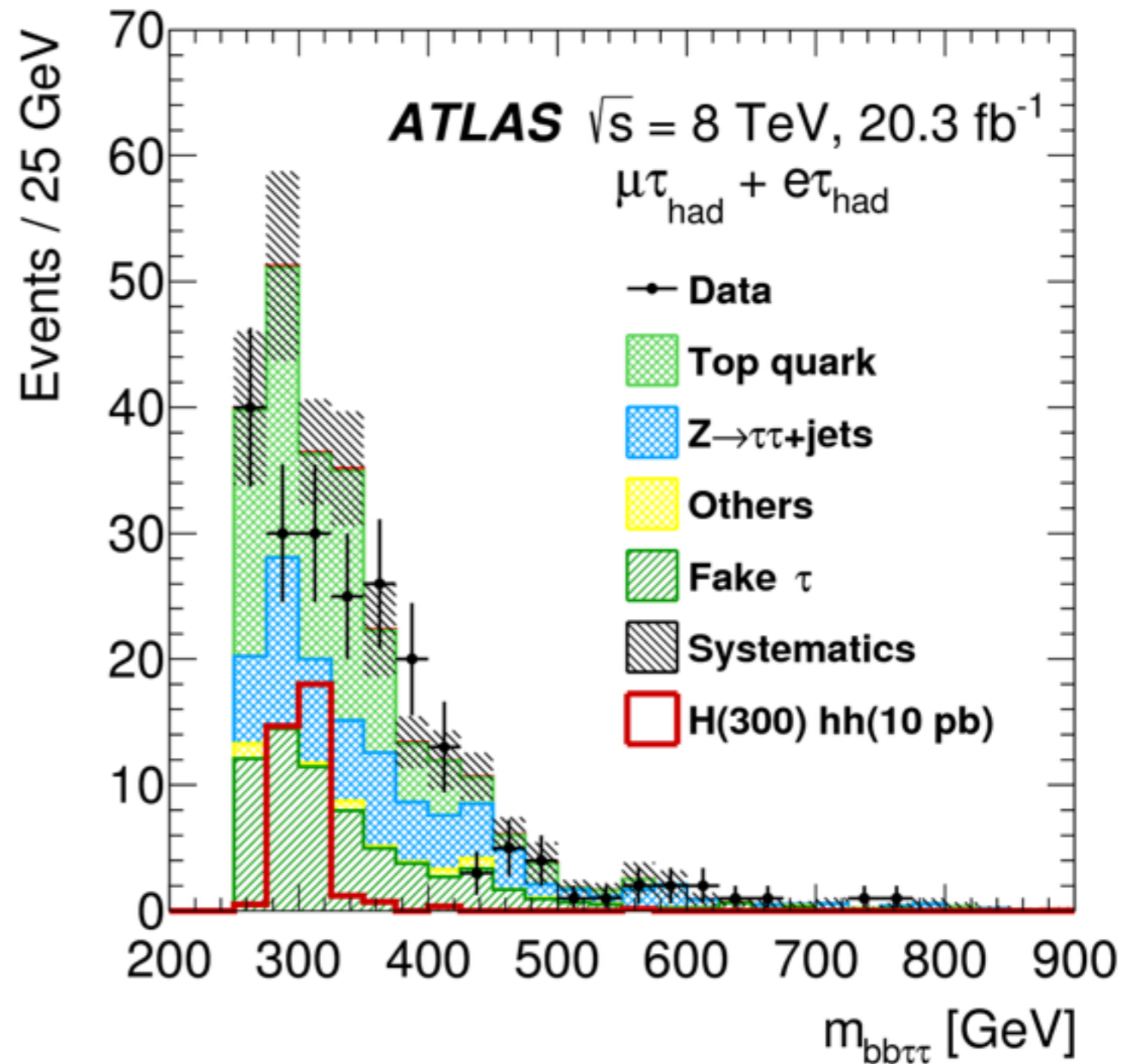
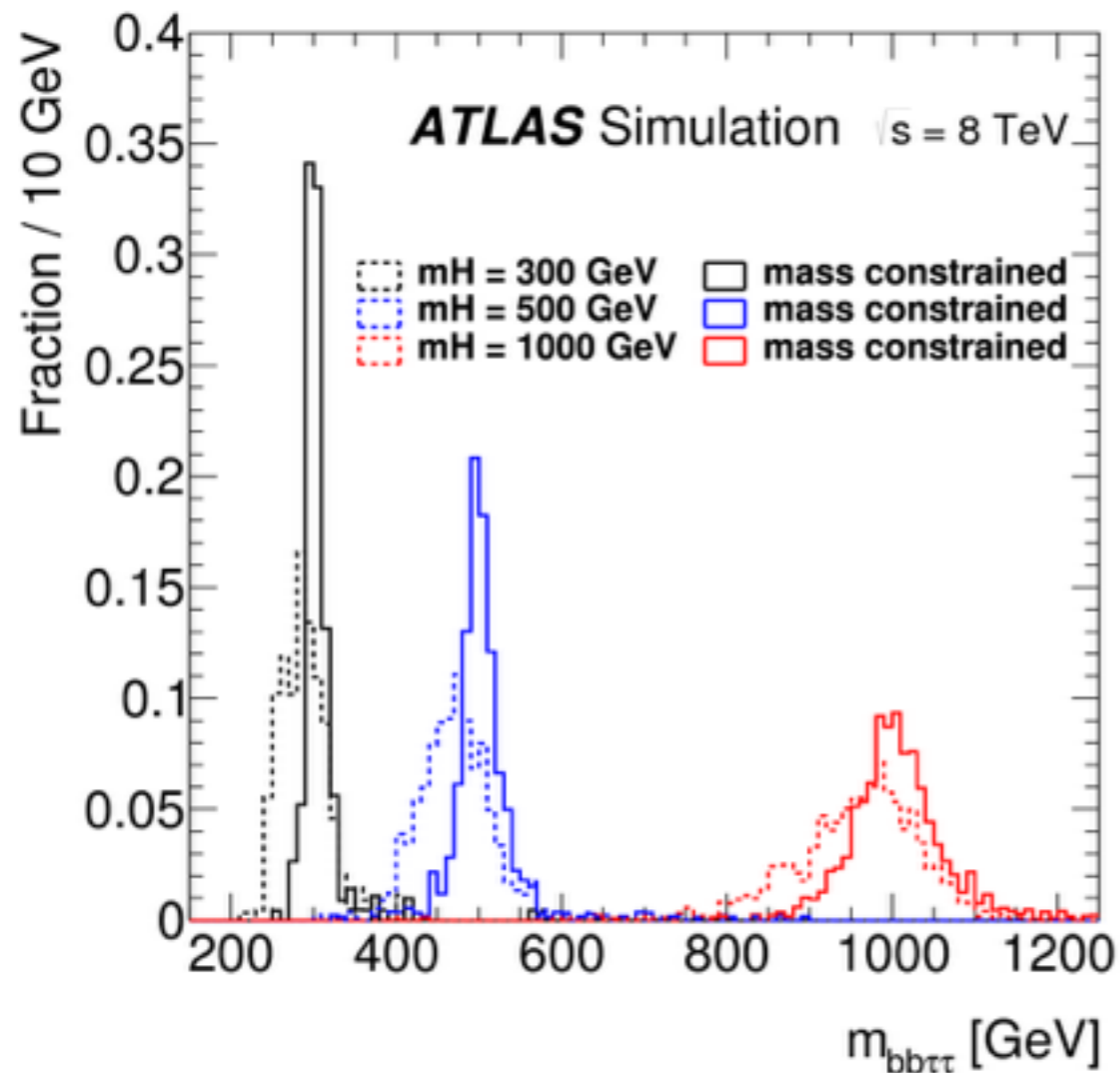
	Cross-section (upper limit)
Expected	1.3 pb
Observed	1.6 pb

	Cross-section (upper limit) relative to SM
Expected	130
Observed	160

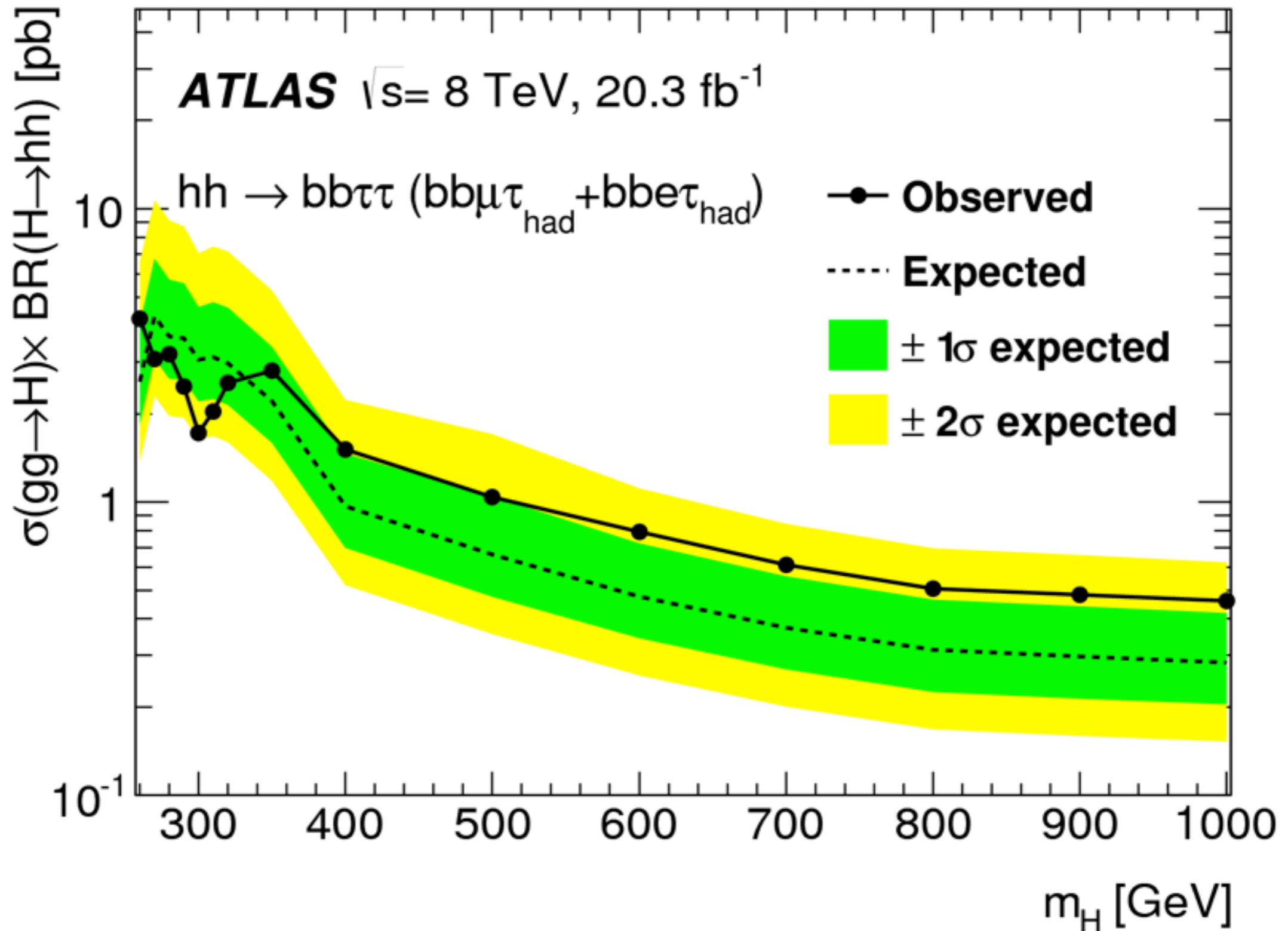


Resonant Analysis

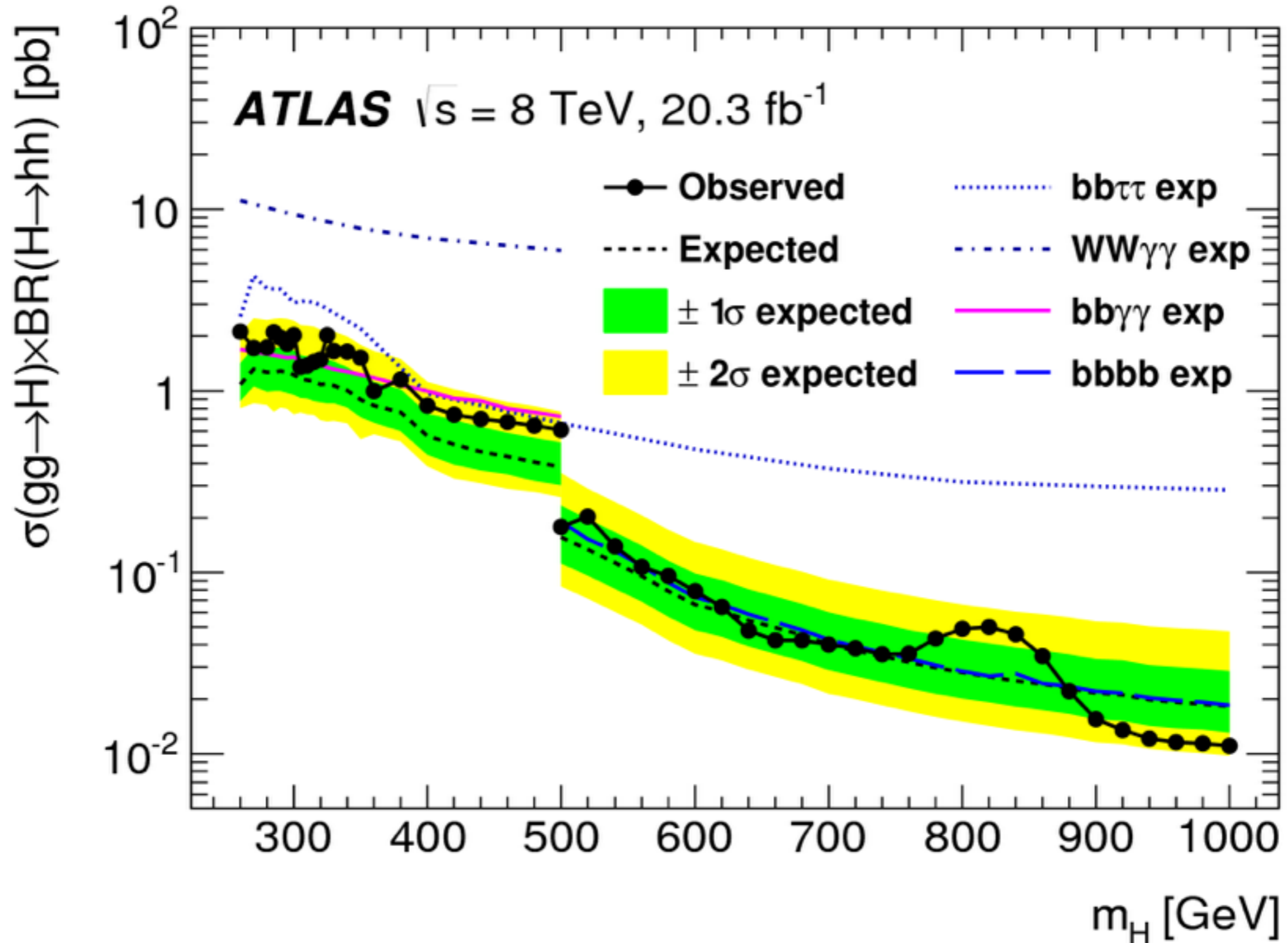
- Apply scale factors m_h/m_{bb} and $m_h/m_{\tau\tau}$ to 4-momenta of bb and $\tau\tau$ systems.
 - ▶ $m_h = 125$ GeV (SM Higgs)
- Improves mass resolution of heavy resonances.



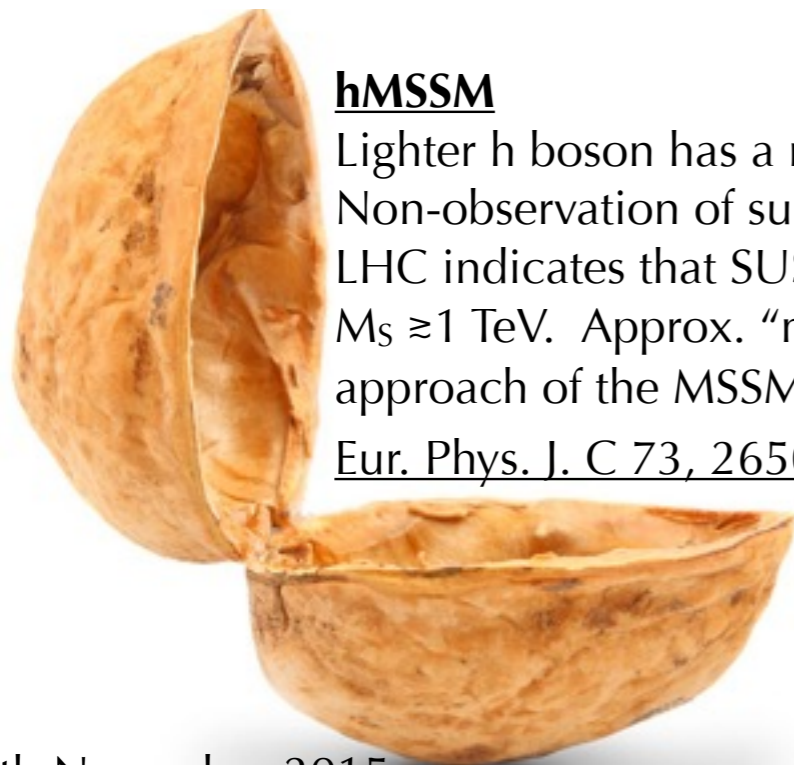
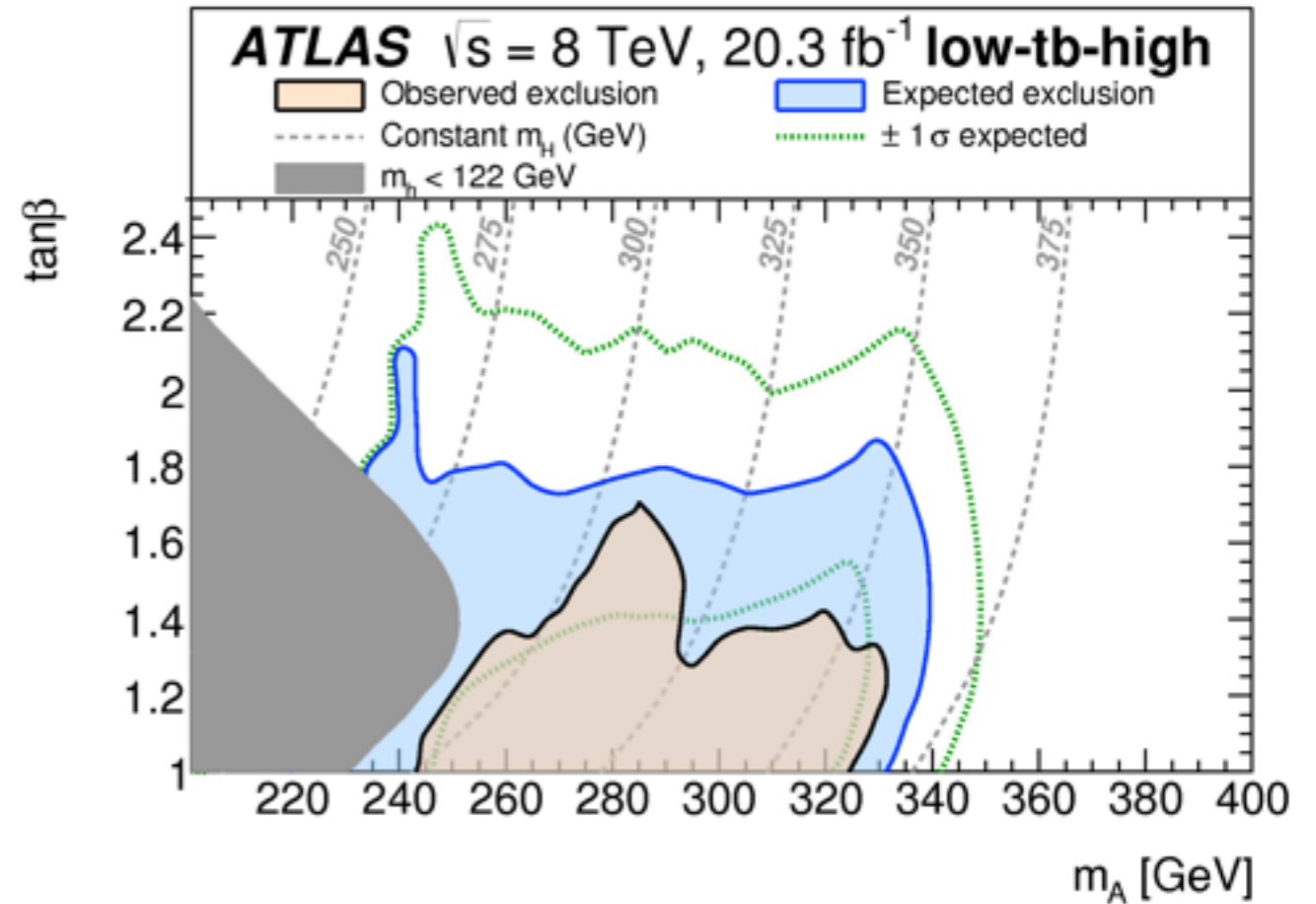
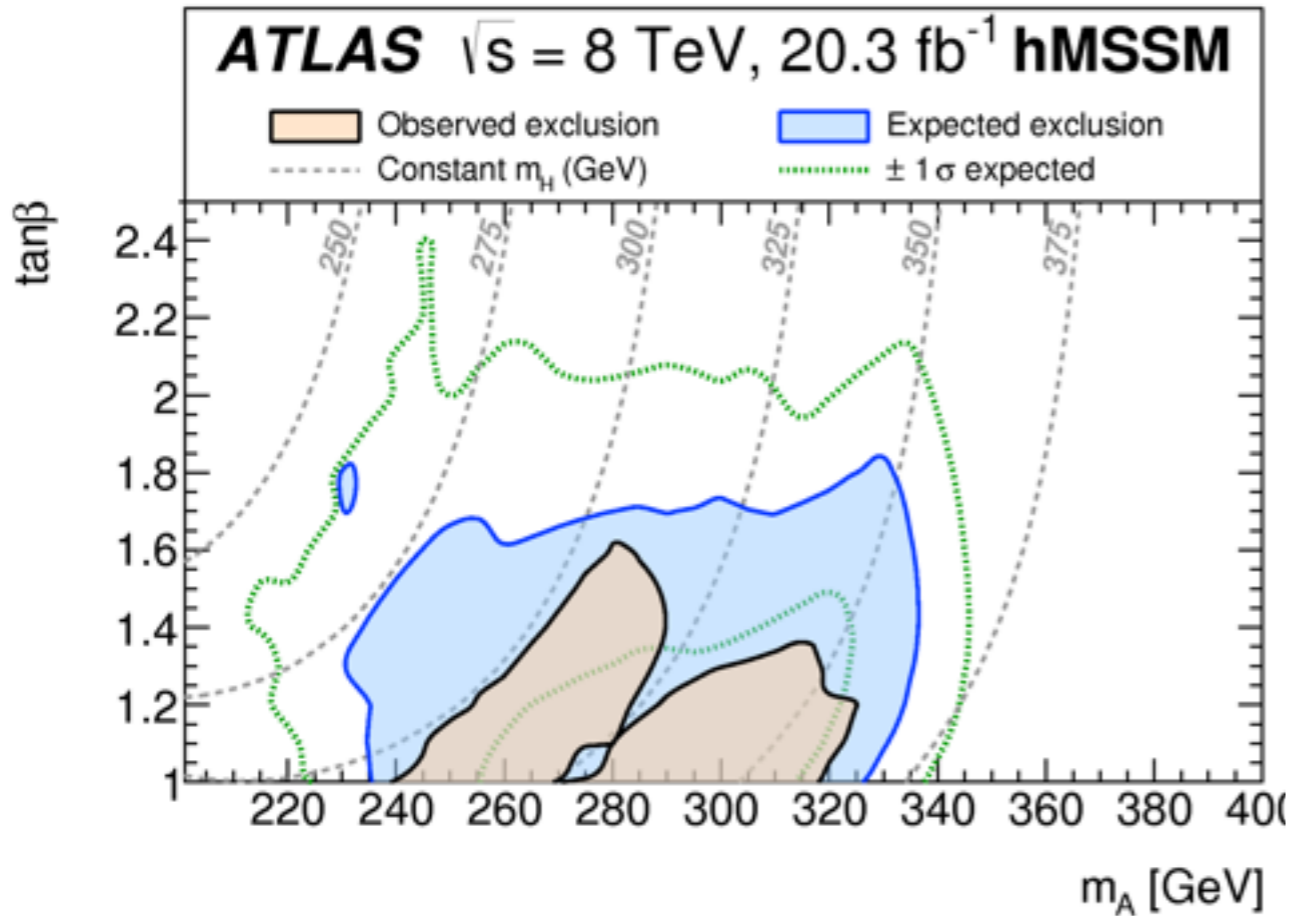
Resonant Analysis



Combination With Other Channels



Constraints on 2HDM Models



hMSSM

Lighter h boson has a mass of 125 GeV. Non-observation of superparticles at the LHC indicates that SUSY-breaking scale $M_s \geq 1 \text{ TeV}$. Approx. "model-independent" approach of the MSSM Higgs sector.

Eur. Phys. J. C 73, 2650 (2013)



low-tb-high:

Lighter h boson has a mass of 125 GeV. Preferred region is low $\tan\beta$ and heavy SUSY.

LHCHSWG-2015-002

RUN 2

NEXT EXIT



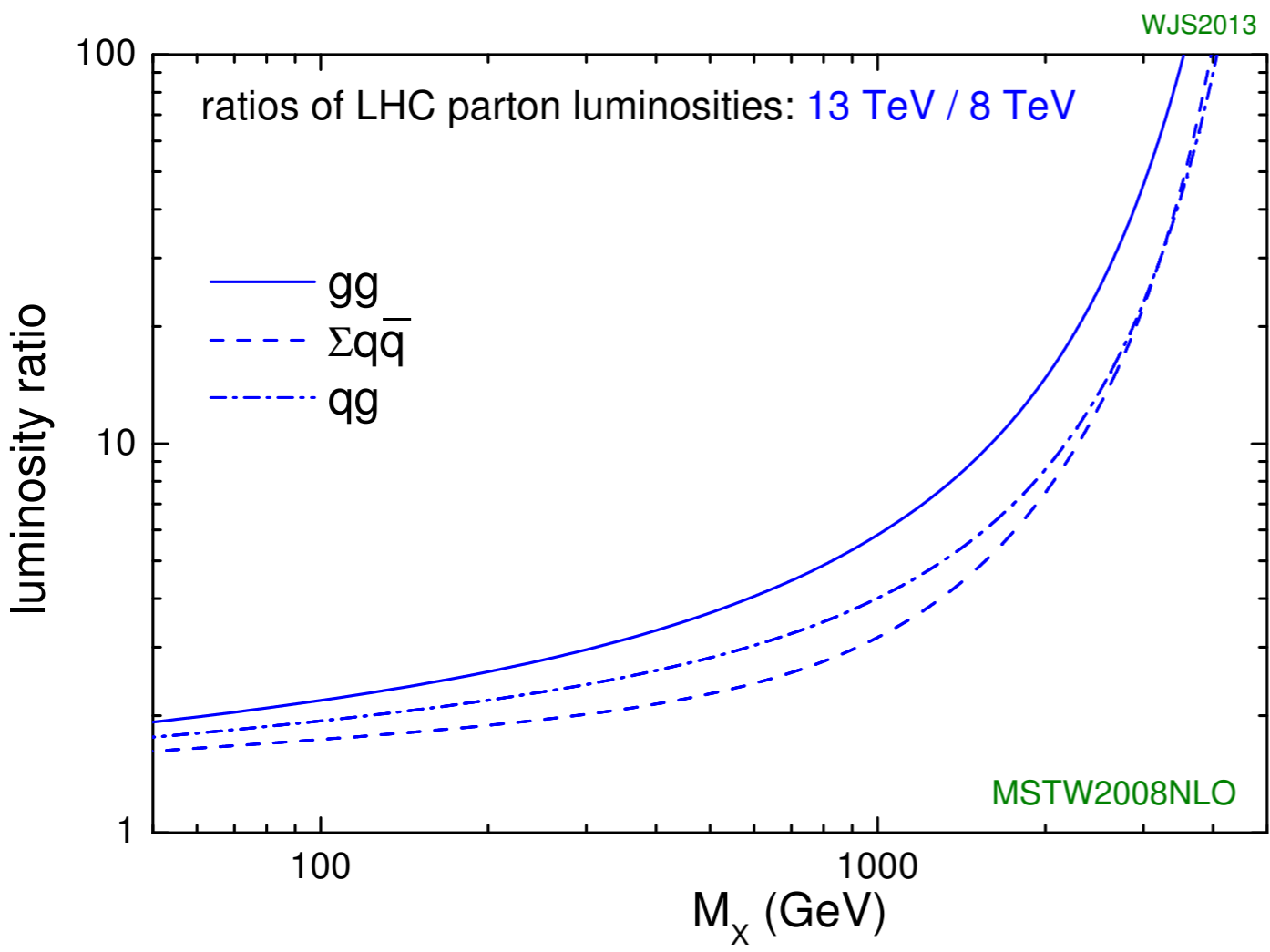
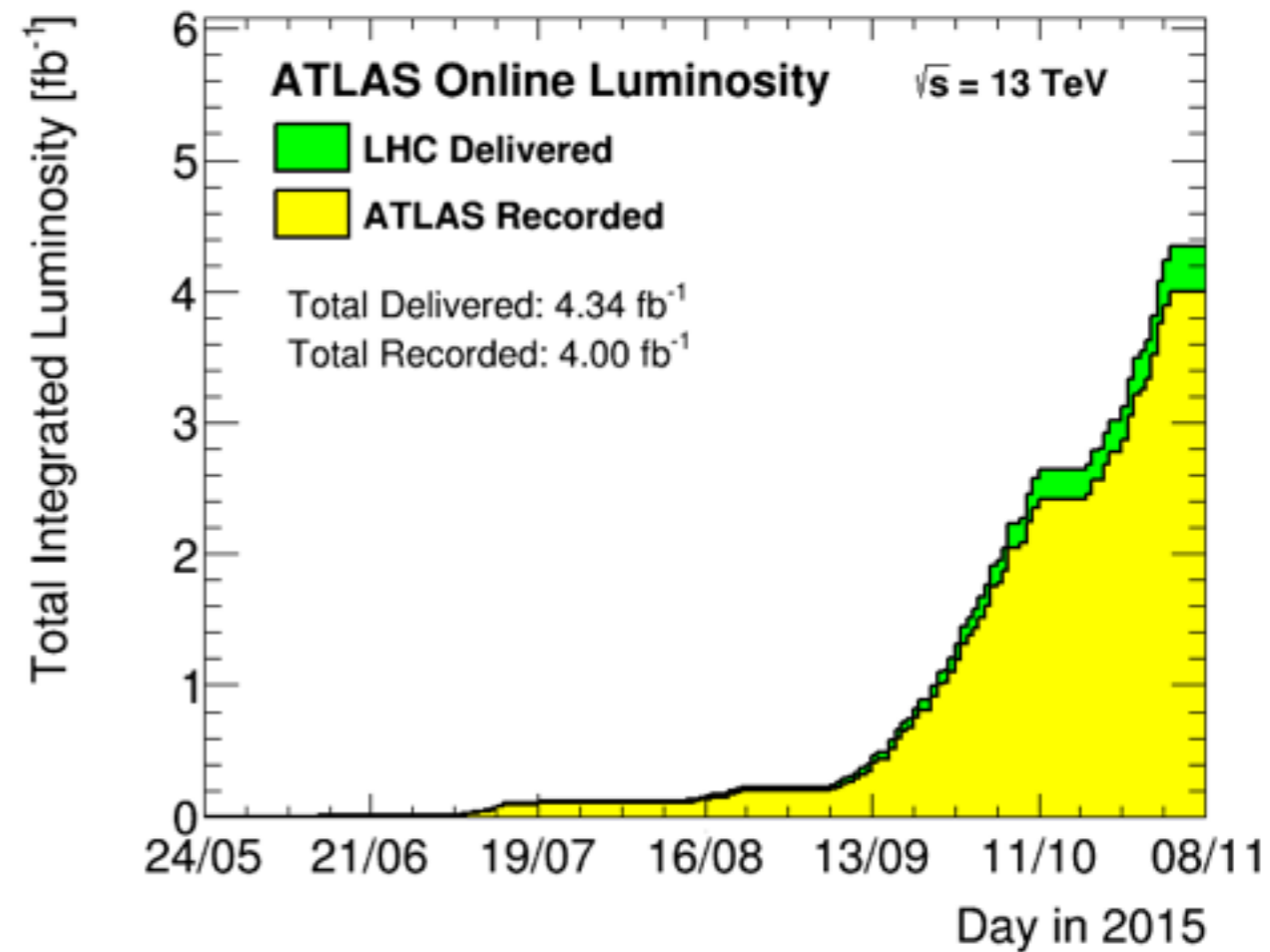
Expected Improvements For Run 2

LHC Run 2:

2015-2018
 $\sqrt{s} = 13 \text{ TeV}$
 100 fb^{-1}
 $\mu = 50$

➔

Access processes
with smaller cross-
sections and/or
higher mass



2015 Summary:

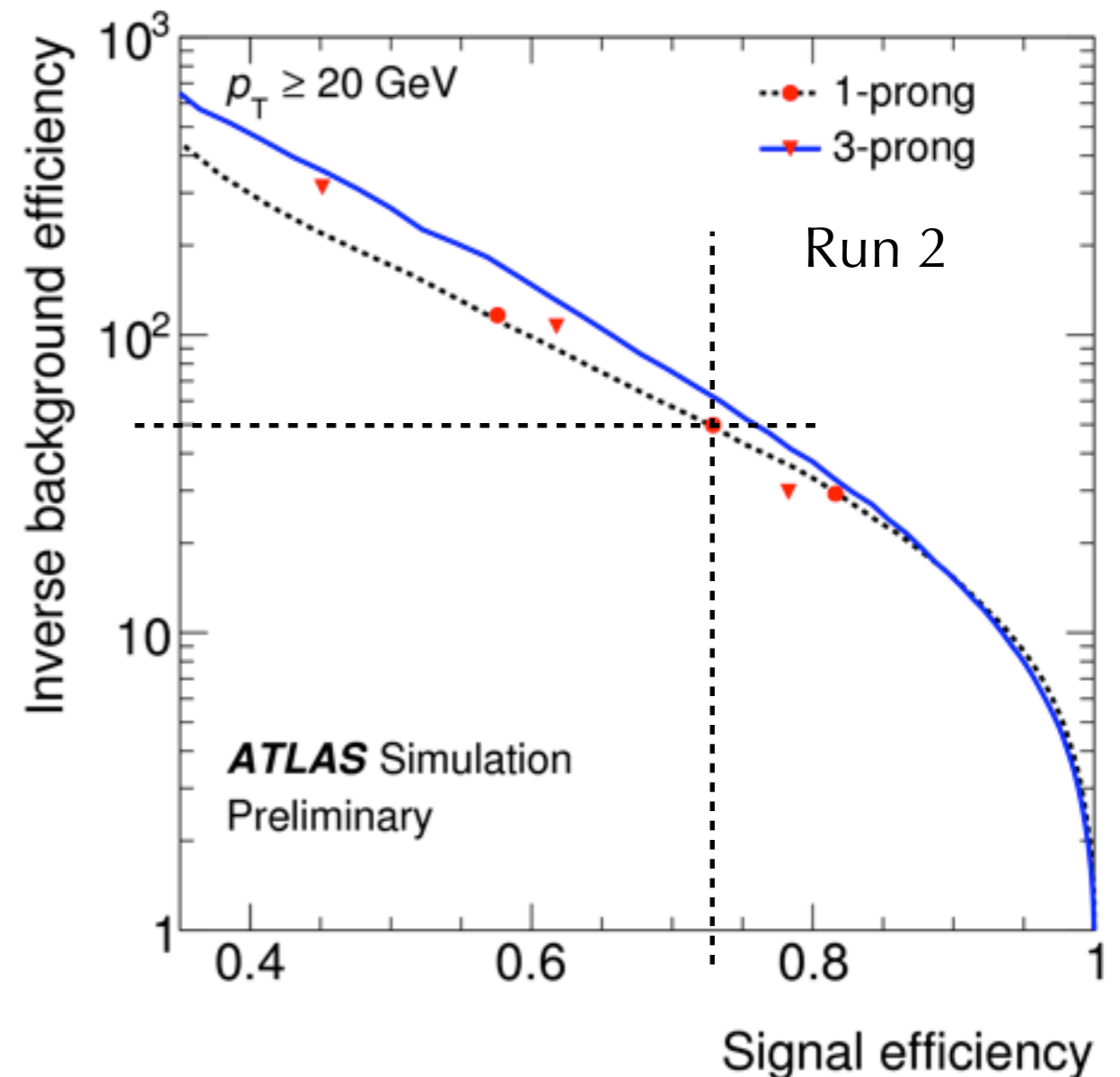
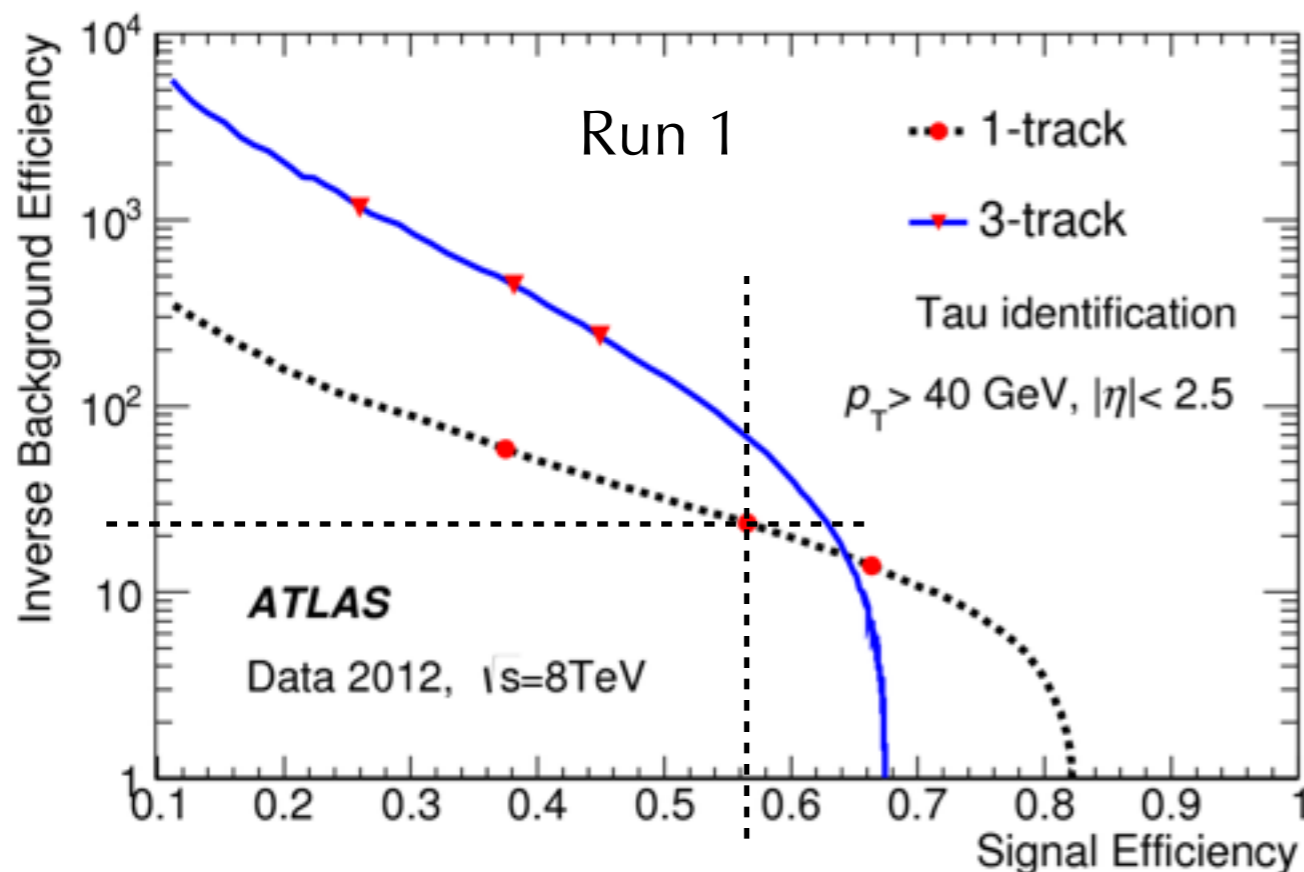
- 4 fb^{-1} collected (3.34 fb^{-1} after data quality requirements).
- Expect to have slightly better sensitivity to $X \rightarrow hh \rightarrow bb\tau\tau$ process than in Run 1 with the 2015 dataset.

Expected Improvements For Run 2

- Include fully hadronic tau-tau decay channel.
 - ▶ Similar sensitivity to lepton-hadron channel.
 - ▶ Fully leptonic channel adds very little, but can be useful as a cross-check.
- Further analysis optimisation.
- Improved object identification for Run 2.

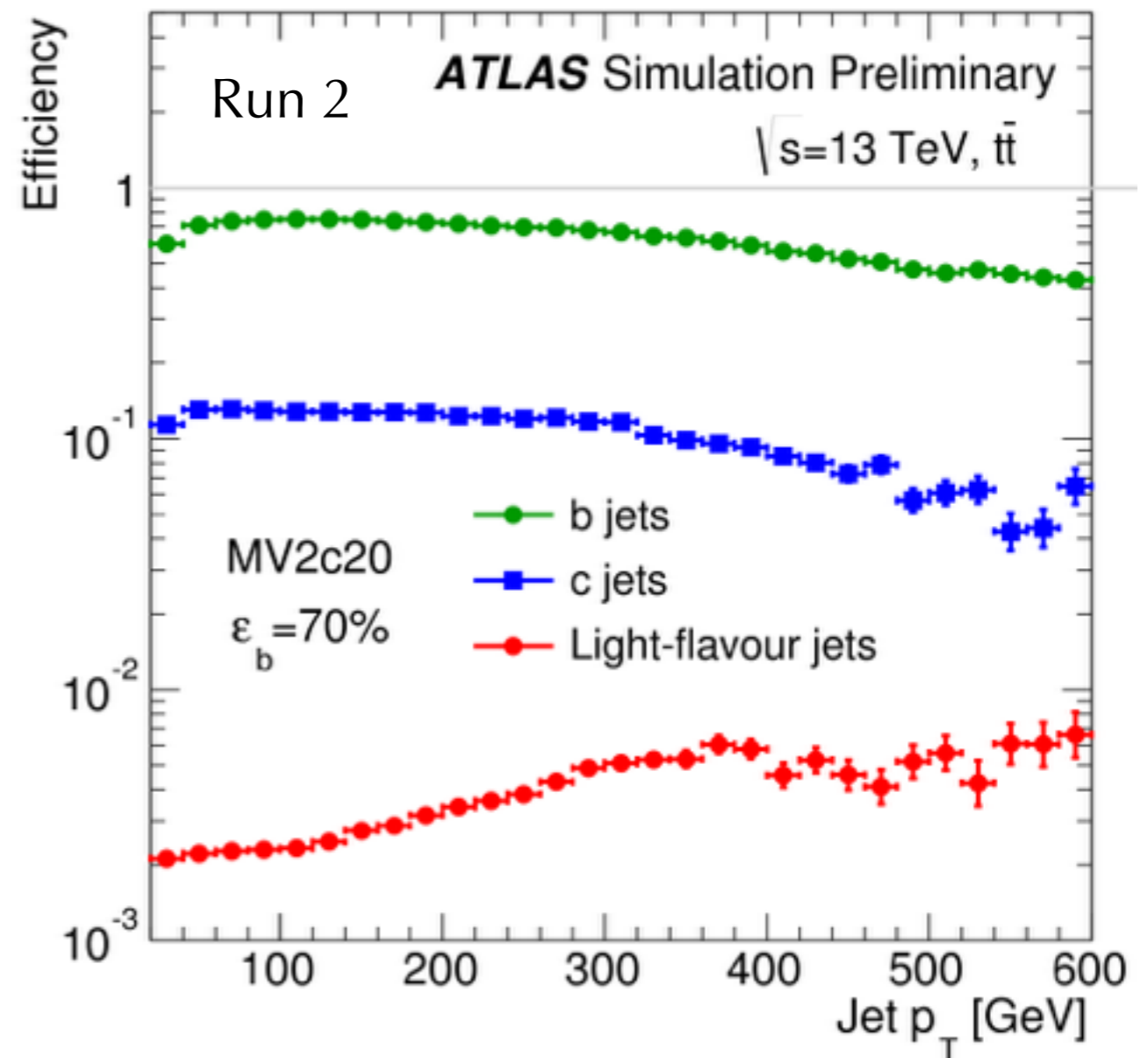
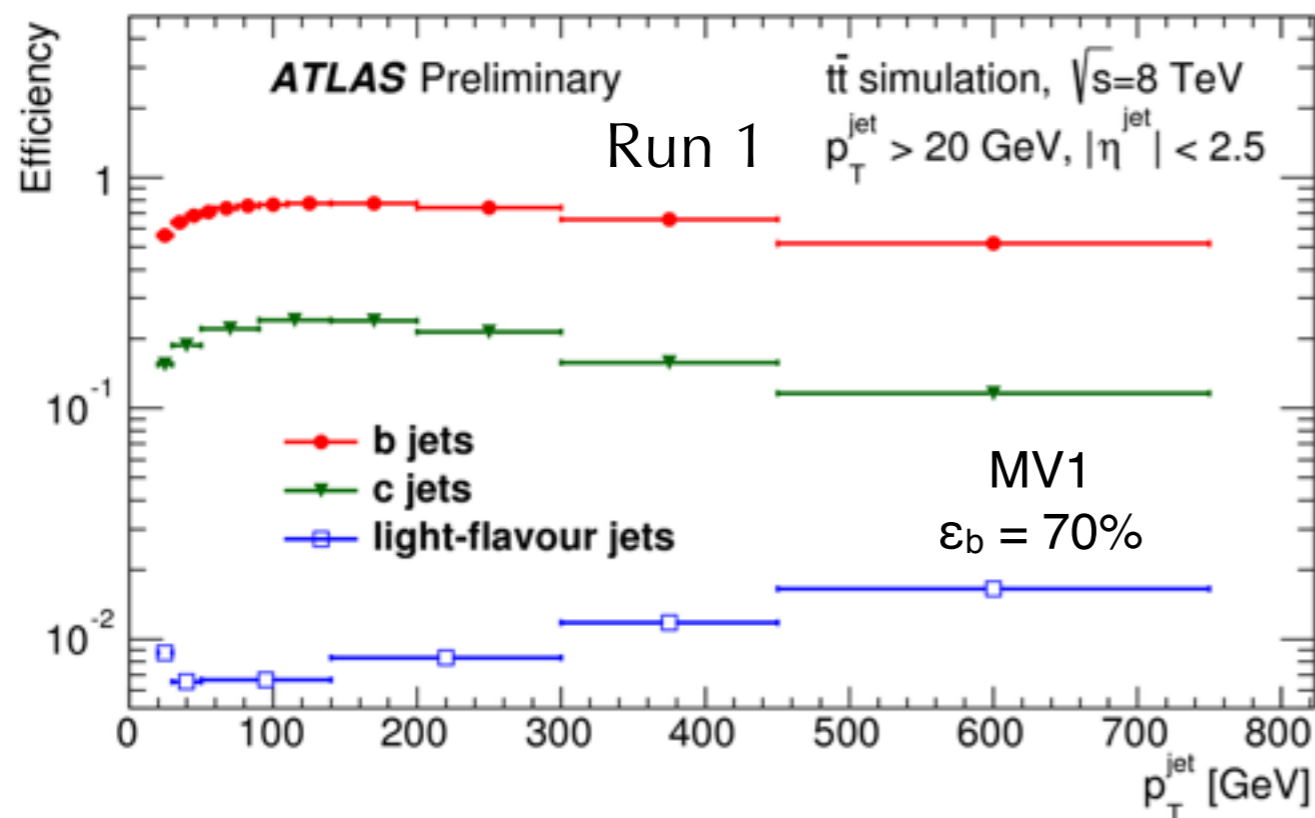
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 - Taus



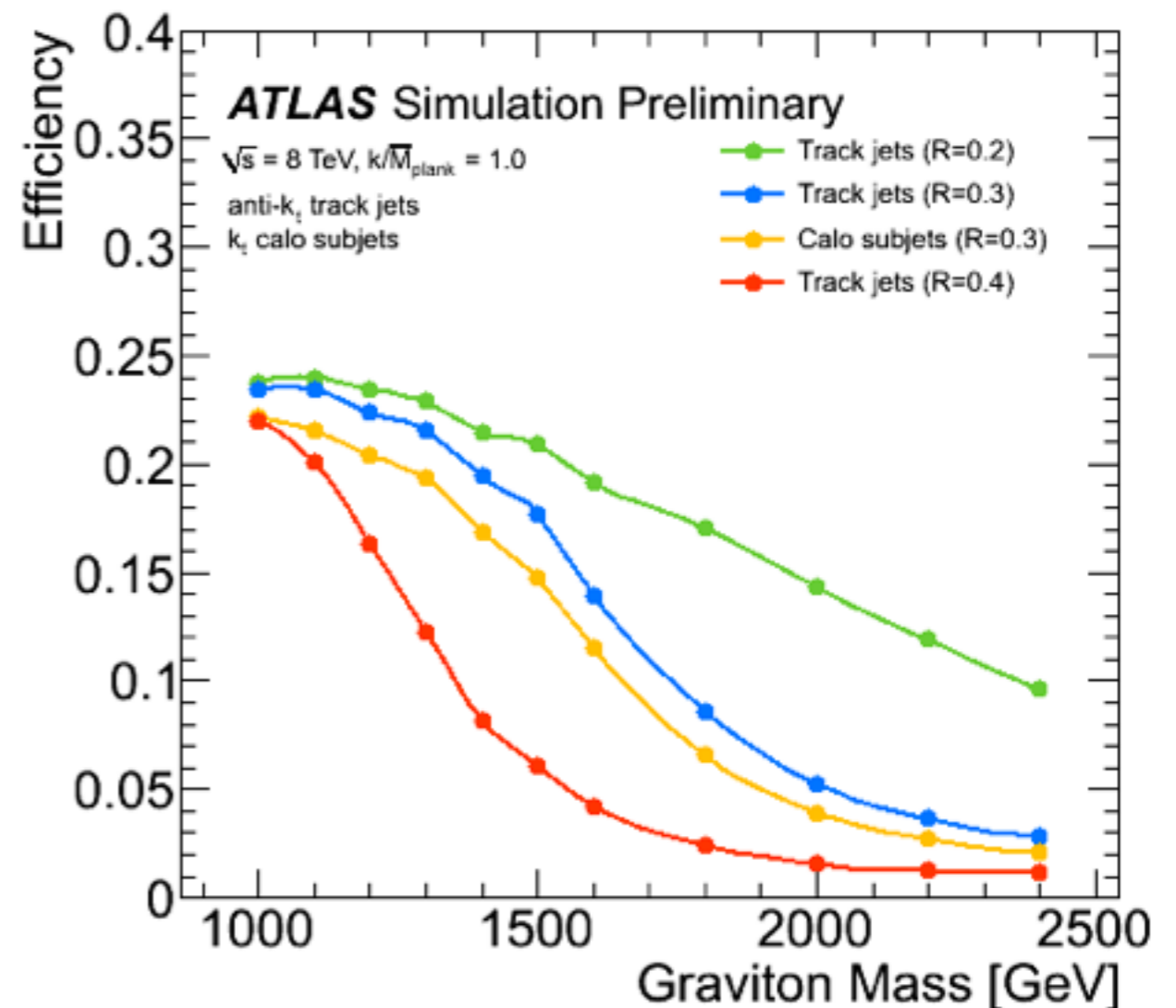
Expected Improvements For Run 2

- Include fully hadronic tau-tau decay channel.
 - Similar sensitivity to lepton-hadron channel.
 - Fully leptonic channel adds very little, but can be useful as a cross-check.
- Further analysis optimisation.
- Improved object identification for Run 2.
 - Taus
 - b-jets



Expected Improvements For Run 2

- Dedicated analysis using sub-structure techniques to reconstruct boosted tau and b-jet pairs will follow later in 2016.
- Current analysis has a natural end-point of ~ 1 TeV, where $\Delta R(\tau, \tau) \sim 0.4$ and tau ID fails.
 - ▶ Normal tau ID relies on an isolation annulus ($0.2 < \Delta R < 0.4$) so fails if the taus are too close together.
 - ▶ Dedicated boosted tau-pair finding algorithm to recover these events and extend mass reach of the analysis.
- b-tagging performance also degrades as a function of $\Delta R(b, b)$.
 - ▶ Less of a 'cliff' than tau ID.
 - ▶ Dedicated tagger for finding boosted pairs of b-jets developed.
 - ▶ $HH \rightarrow bbbb$ analysis shows significant gains when using it.

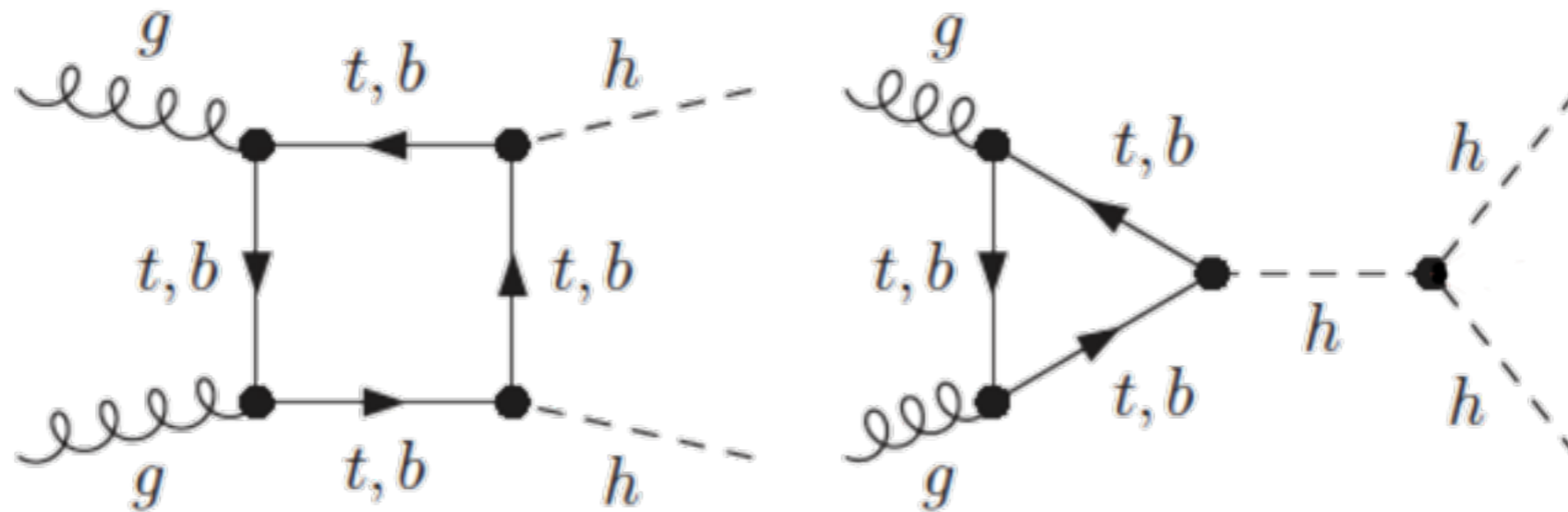


A green highway sign with white text, mounted on a metal structure. The sign is rectangular with rounded corners and a white border. The background is a clear blue sky with a yellowish glow at the bottom.

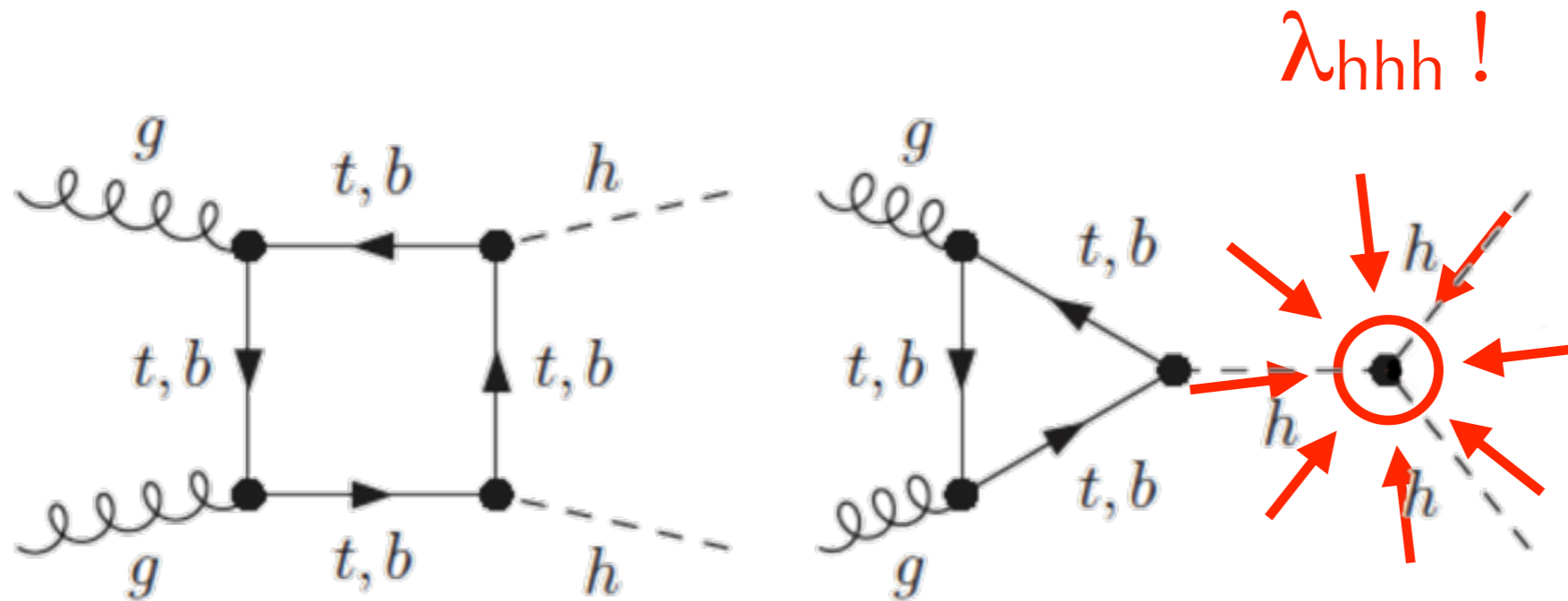
HL-LHC

QUITE A LONG WAY OFF
(2023-2035-ISH)

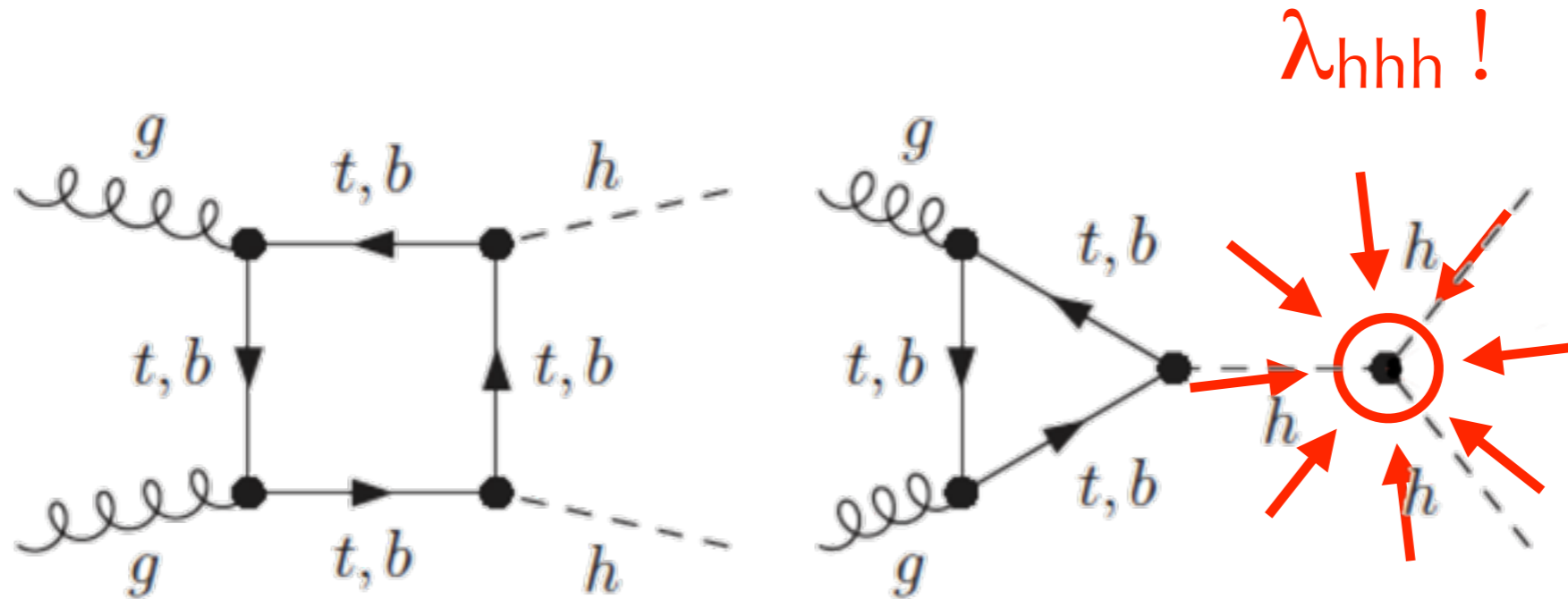
SM Higgs Pair Production



SM Higgs Pair Production

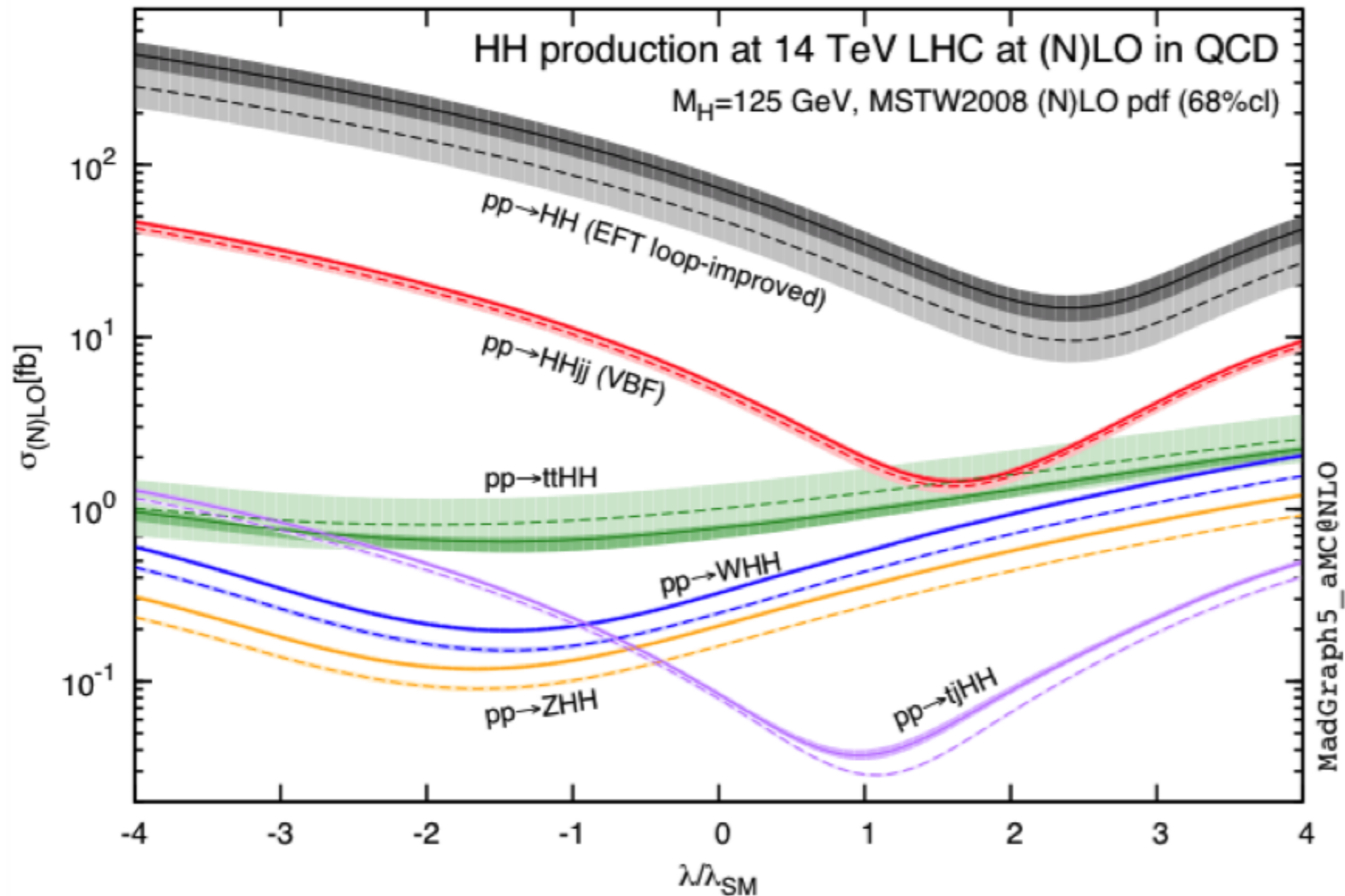


SM Higgs Pair Production



- Higgs self-coupling crucial check of EWSB mechanism.
- Possibly *the* most challenging measurement at the LHC!
- Direct measurement of the Higgs trilinear self-coupling (λ_{HHH}) can be made by studying Higgs pair-production.
- Need to dis-entangle top box-diagram and diagram containing the HHH vertex...
- Destructive interference with diagrams not containing the HHH vertex.
 - Box diagram dominates in boosted events.
 - Absolutely crucial to push down to lower p_T 's in order to access λ_{HHH} .

SM Higgs: Pair Production



- For $\lambda_{HHH}/\lambda_{HHH}^{SM} = 0/1/2$, cross-section = 71/34/16 fb.
- With 3000 fb^{-1} a $\sim 3\sigma$ combined measurement by ATLAS+CMS should be possible.

SM $HH \rightarrow bb\tau\tau$ Sensitivity Study

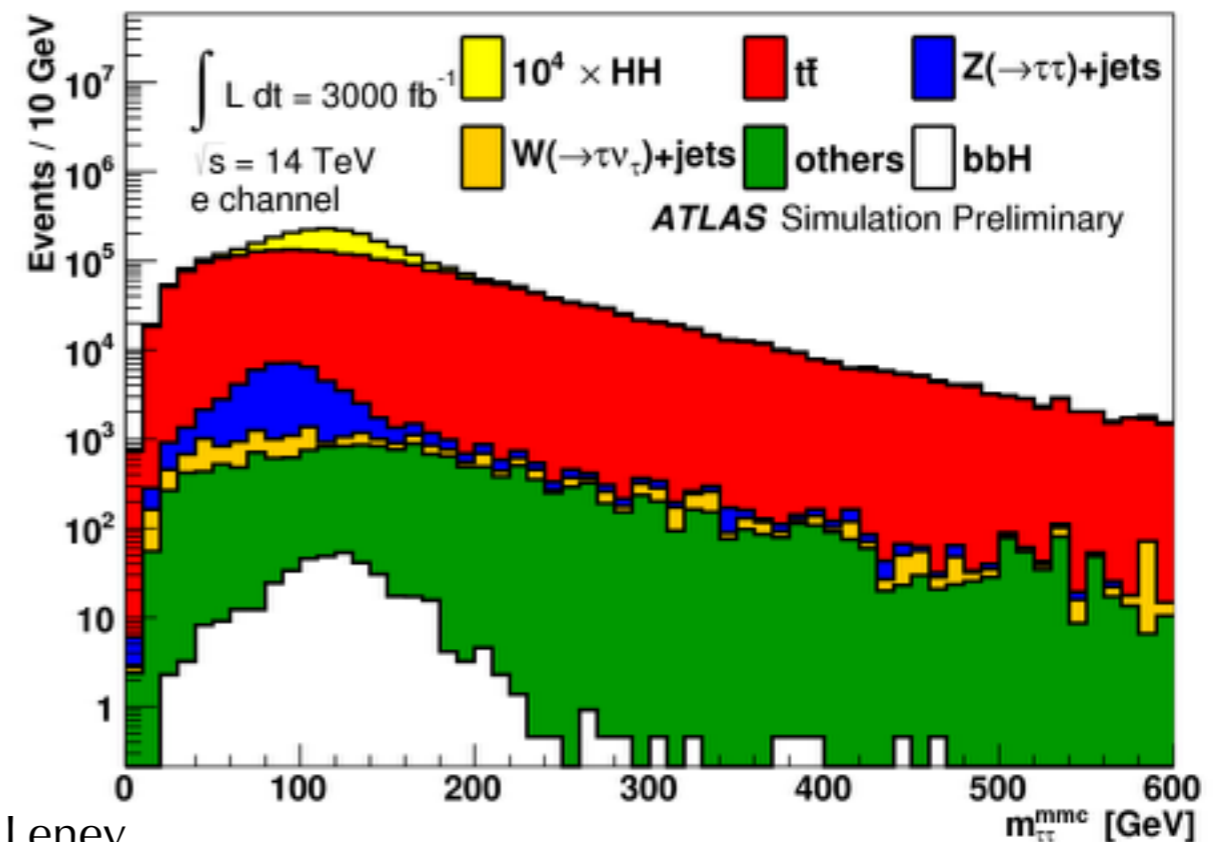
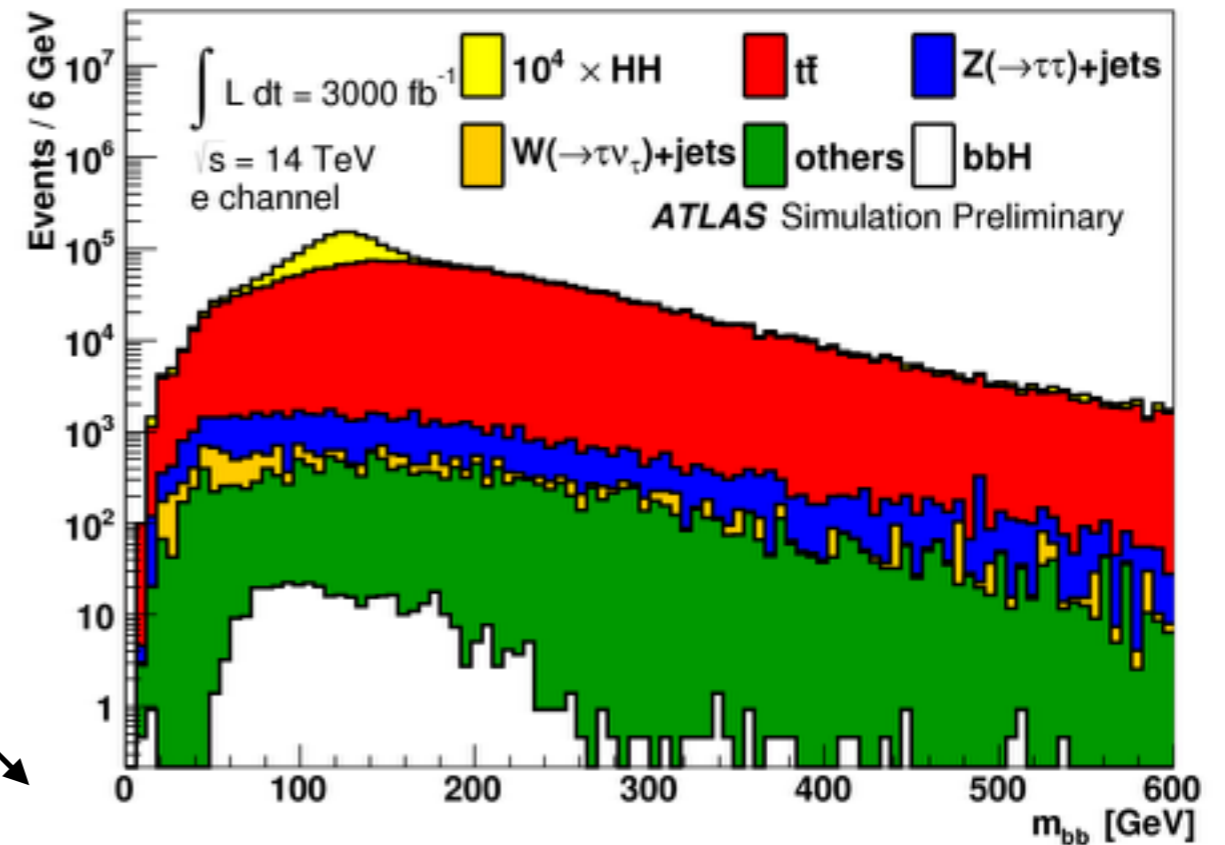
ATL-PHYS-PUB-2015-046

- Truth-level, 'cut-and-count' study for first estimate of sensitivity.
 - ▶ Detector response parameterised.
 - ▶ More sophisticated analyses (using MVAs) will be necessary.
- Assume 3000 fb^{-1} data, $\sqrt{s} = 14 \text{ TeV}$, and 140 proton-proton collisions per bunch-crossing.
- All tau-tau decay modes considered:
 - ▶ Fully leptonic ($\tau_\ell\tau_\ell$)
 - ▶ Semi-leptonic ($\tau_\ell\tau_h$)
 - ▶ Fully hadronic ($\tau_h\tau_h$)
- SM Higgs processes that have been negligible backgrounds in the new physics searches so far become more important.
 - ▶ e.g. VH , ttV , ttH
- Higgs bosons produced by λ_{HHH} process have low p_T .
 - ▶ Those produced via top box-diagram are more boosted.
 - ▶ Lower p_T objects harder to separate from multi-jets backgrounds.
 - ▶ Need to find the right balance...

Event Selection

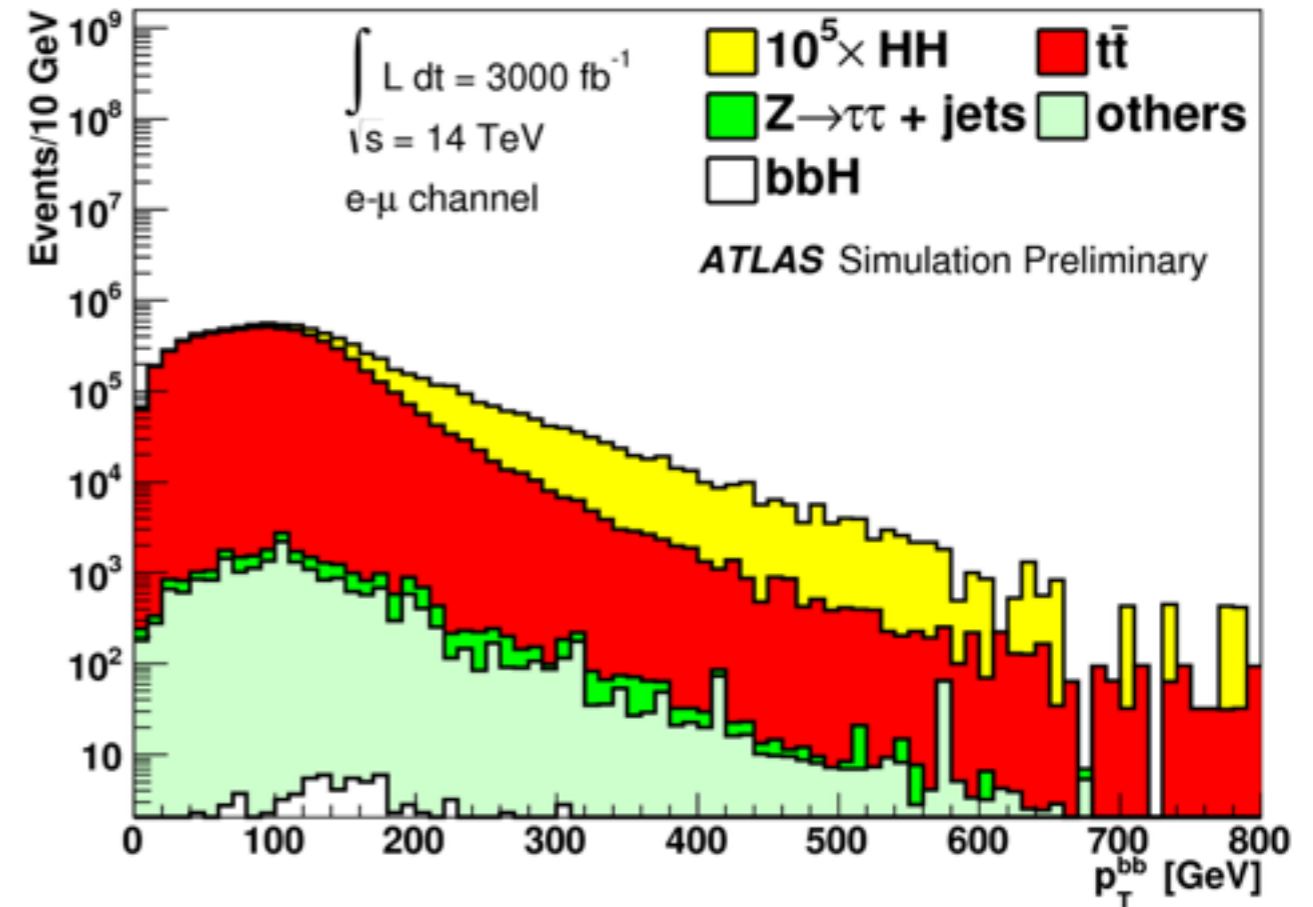
	Example Cut Value*
m_{bb} mass window	$95 < m_{bb} < 145$ GeV
$m_{\tau\tau}$ mass window	$90 < m_{\tau\tau} < 160$ GeV
$p_T(bb)$	$p_T(bb) > 200$ GeV
bb-pair separation	$\Delta R(bb) < 1.0$
Transverse mass	$m_T < 80$ GeV
Stransverse mass	$m_{T2} > 180$ GeV

* Slight variations in exact cut values between channels



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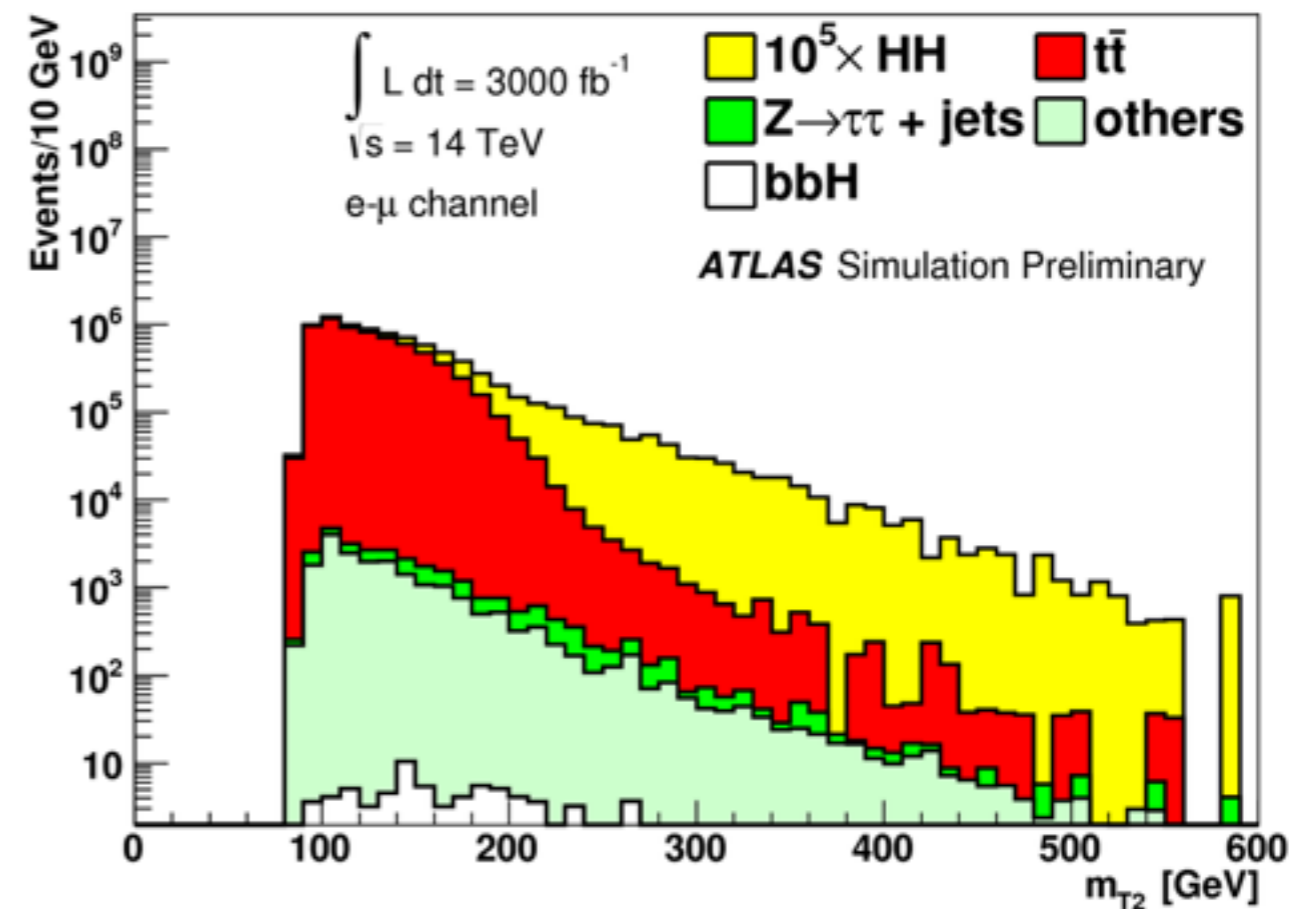
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Stransverse mass: Generalisation of the transverse mass when applied to signatures with two (or more) invisible particles in the final state.

Phys.Lett. B463 (1999) 99–103



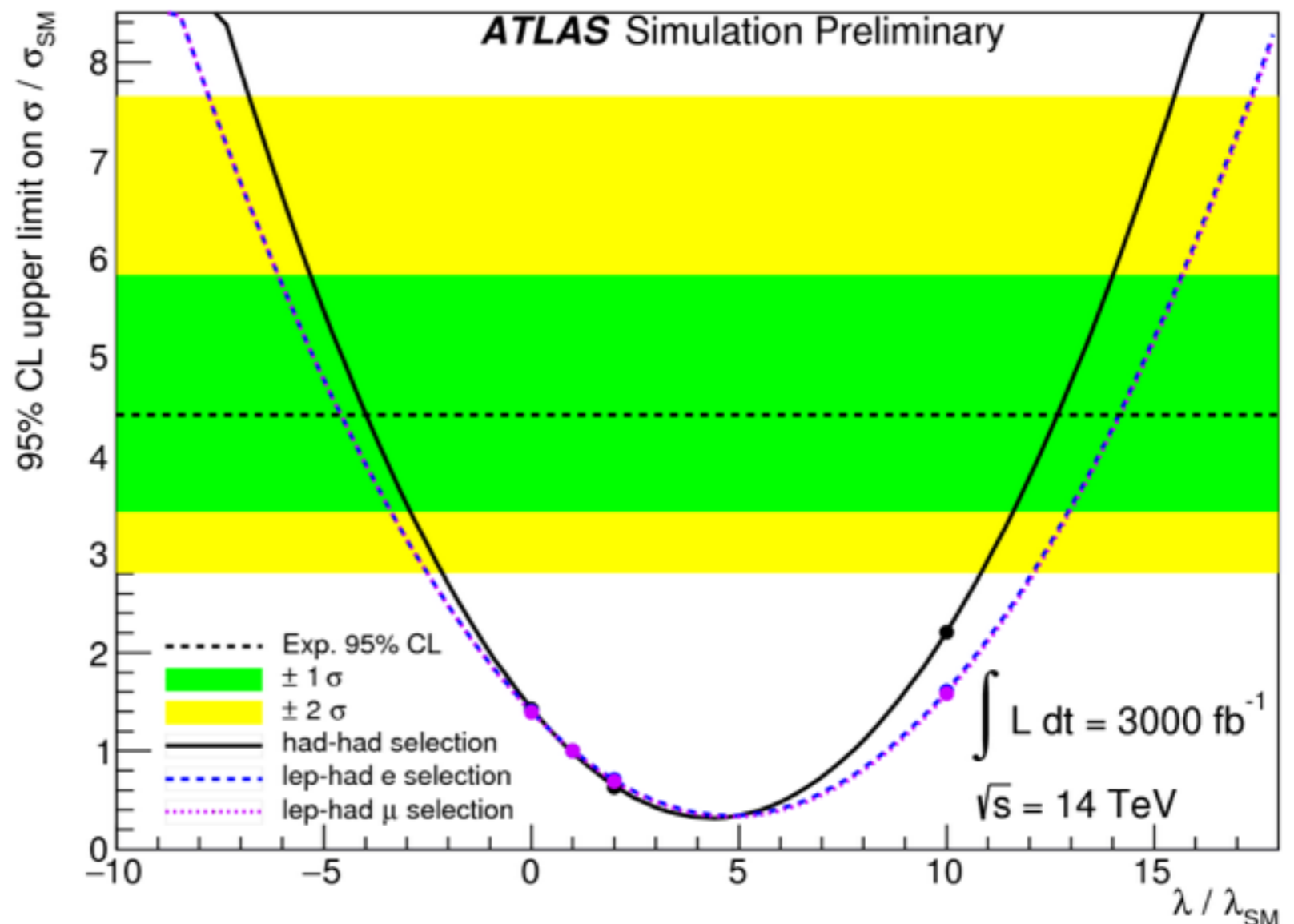
Prospects for Measuring λ_{HHH} at the HL-LHC

Channel	Significance	Combined in channel	Total combined
$e + \text{jets}$	0.31	0.43	0.60
$\mu + \text{jets}$	0.30		
$\mathcal{T}_{\text{had}}\mathcal{T}_{\text{had}}$	0.41	0.41	

$\tau_\ell\tau_\ell$ channels order of magnitude less sensitive than $\tau_\ell\tau_h$ or $\tau_h\tau_h$ channels.

- Use of more sophisticated analysis techniques can improve sensitivity.
- Combination across many channels will be necessary.
- Large correlation between total di-Higgs production cross-section and λ_{HHH} .
 - λ_{HHH} better studied using shape analysis of key observables.

$\sigma / \sigma_{\text{SM}}$ as a function of $\lambda / \lambda_{\text{SM}}$



Summary & Conclusions

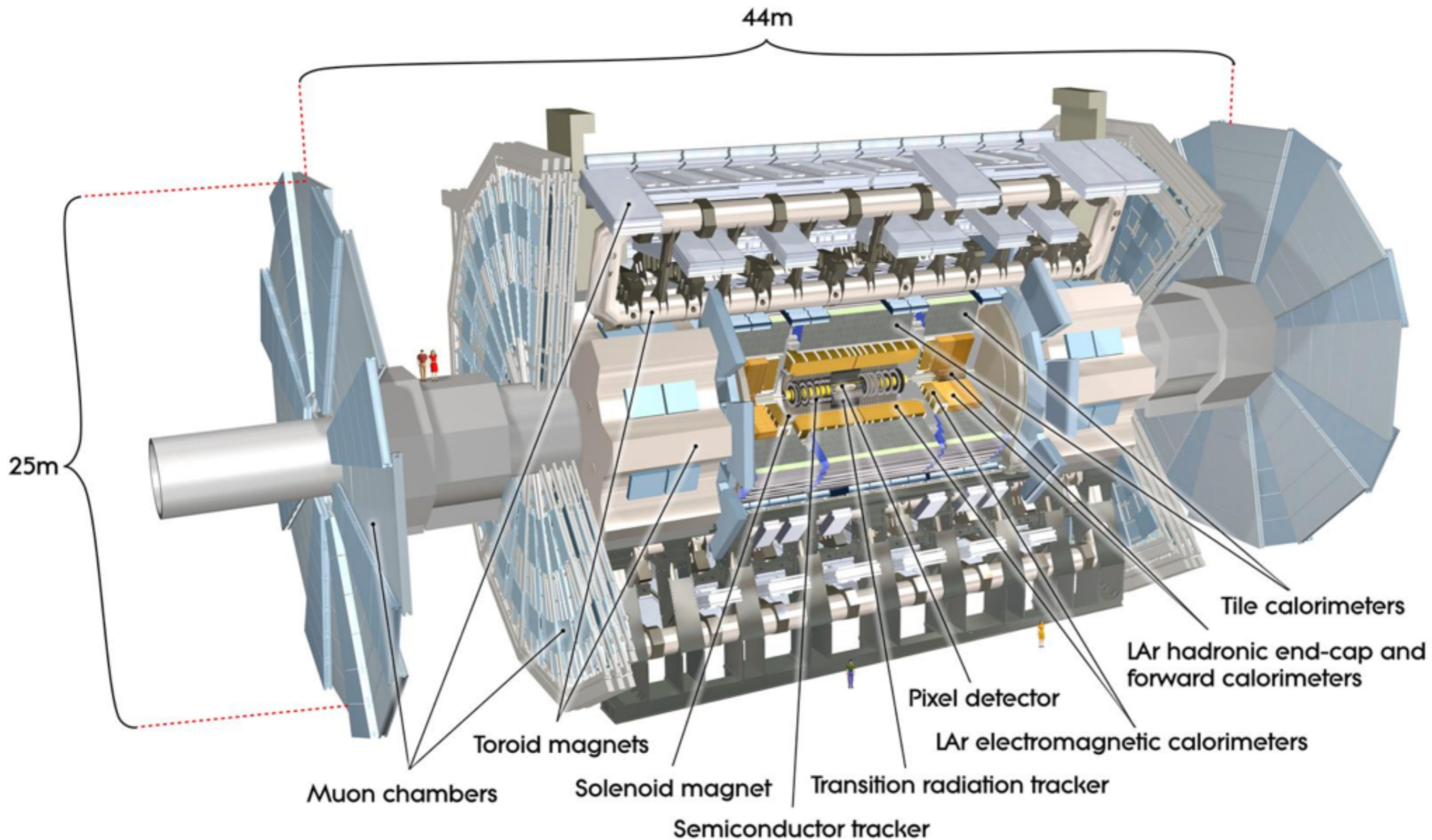
- Di-Higgs physics is a new and exciting avenue for searching for new physics.
 - ▶ Many new physics models predict enhanced rates, either as resonant or non-resonant production.
 - ▶ Also predicted by the SM where it will be a crucial test of the Higgs mechanism.
 - ▶ $bb\tau\tau$ is a promising channel, particularly in the SM/lower mass/non-resonant regime.
- Run 1 searches showed no signal in the $bb\tau\tau$, or any other channel.
 - ▶ Limits set on resonant and non-resonant HH production.
- Will significantly extend this reach during the LHC Run 2.
 - ▶ Even with 2015 data alone should do better than Run 1.
 - ▶ Extra analysis optimisations will improve the sensitivity further, particularly at higher masses.
- The $bb\tau\tau$ channel is a promising one for measuring di-Higgs production and λ_{HHH} at the HL-LHC.



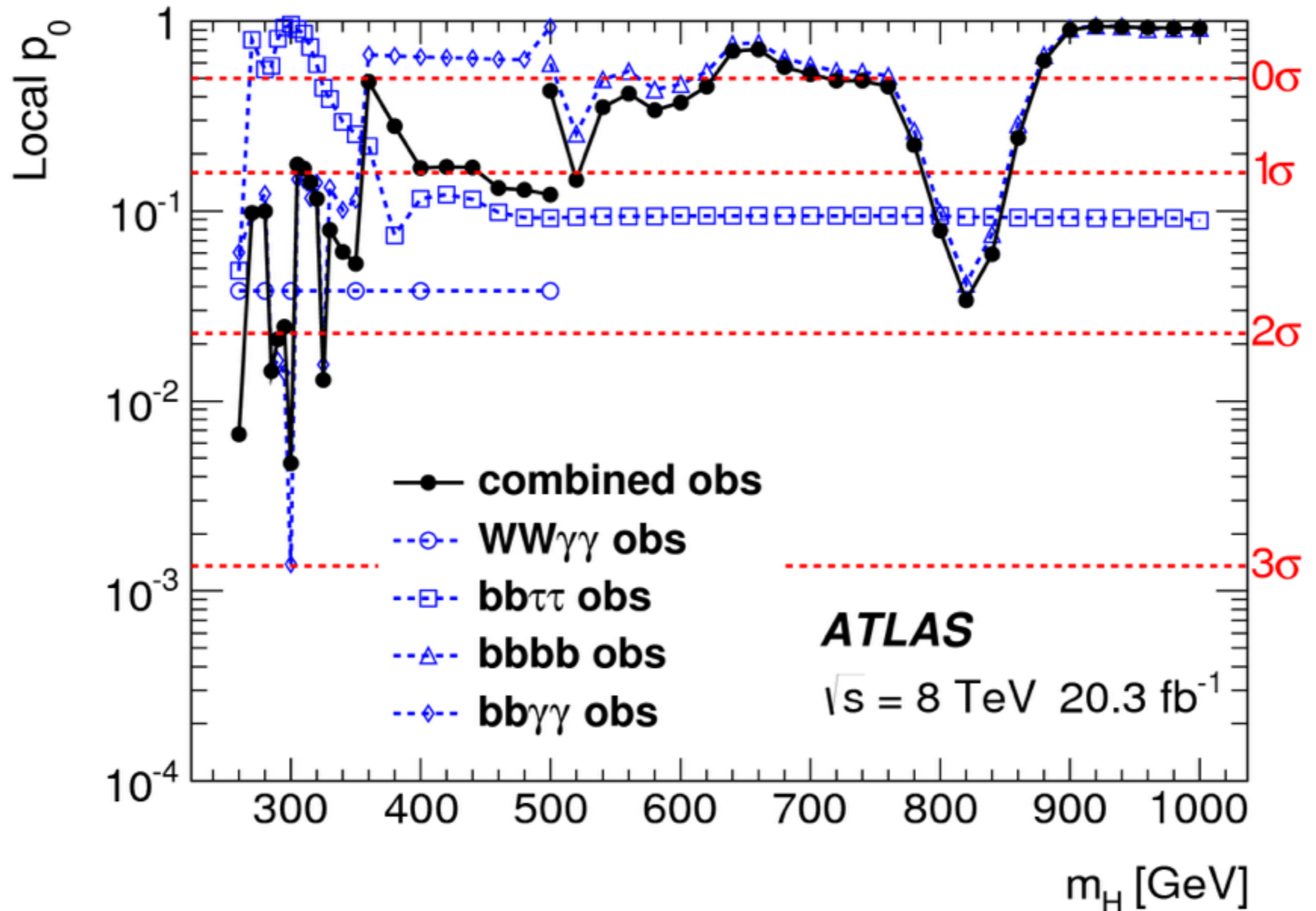
Is there anything beyond the Standard Model?

Back Up

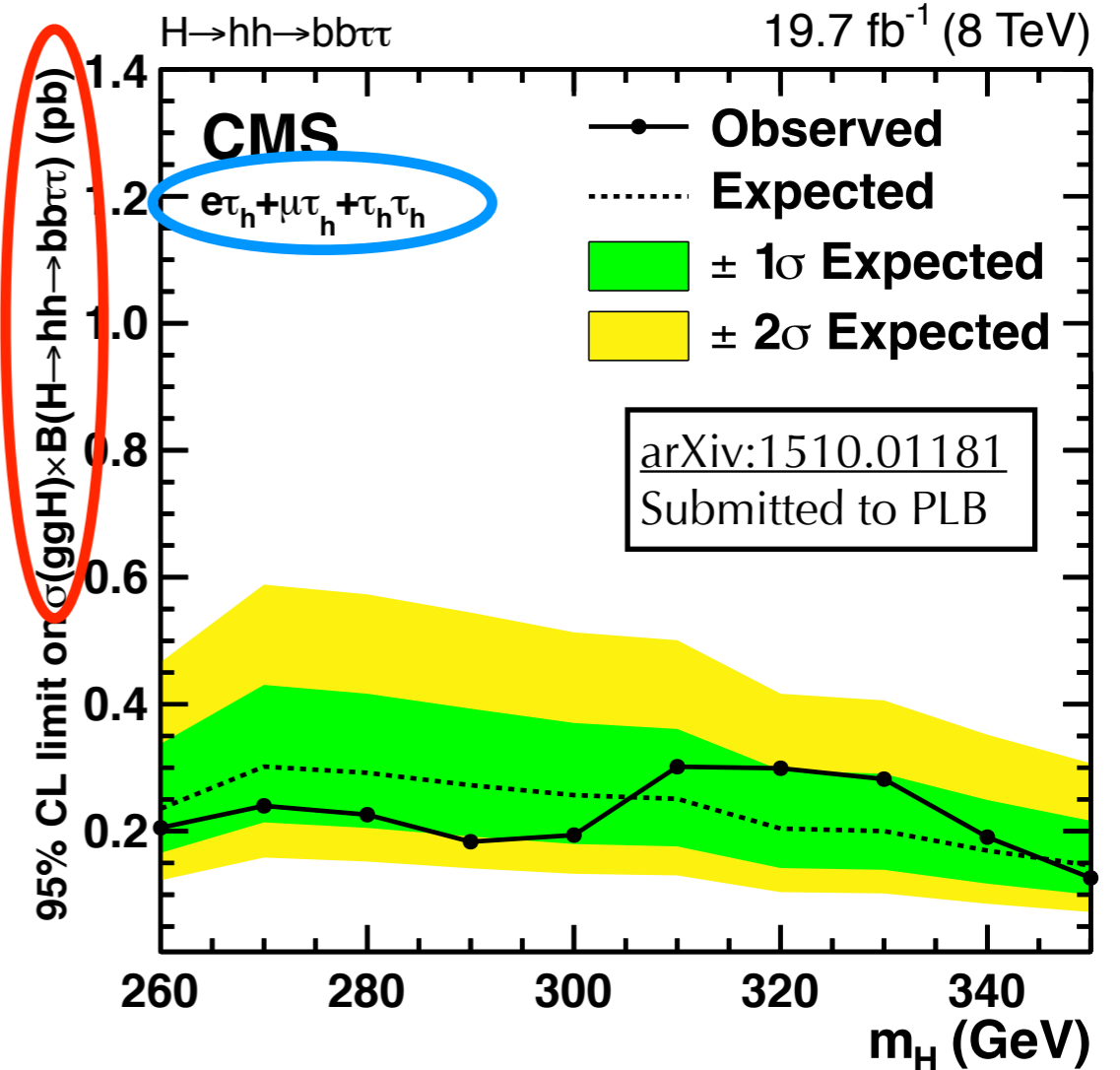
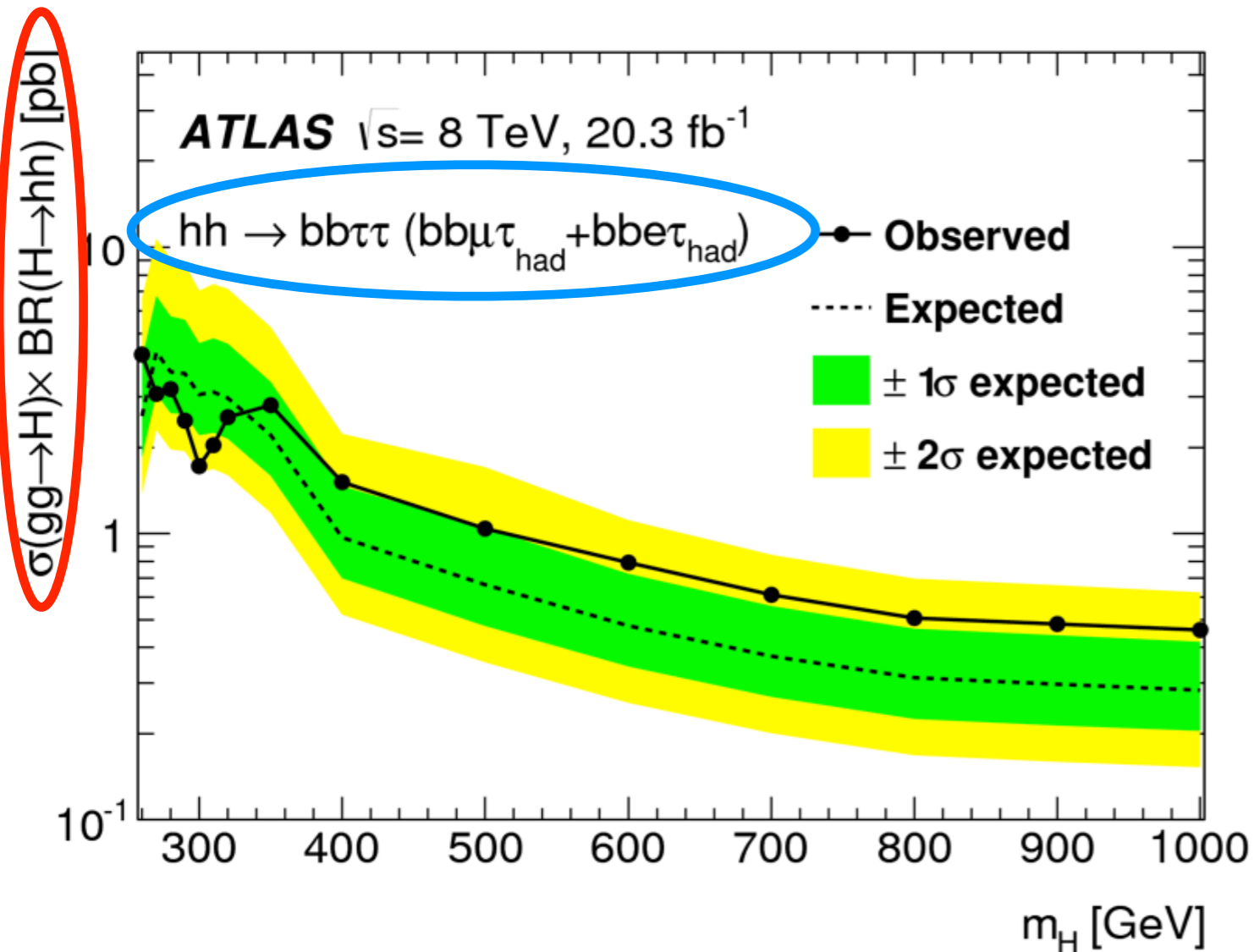
Obligatory ATLAS Detector Slide



Combination With Other Channels



Resonant Analysis - Comparison with CMS

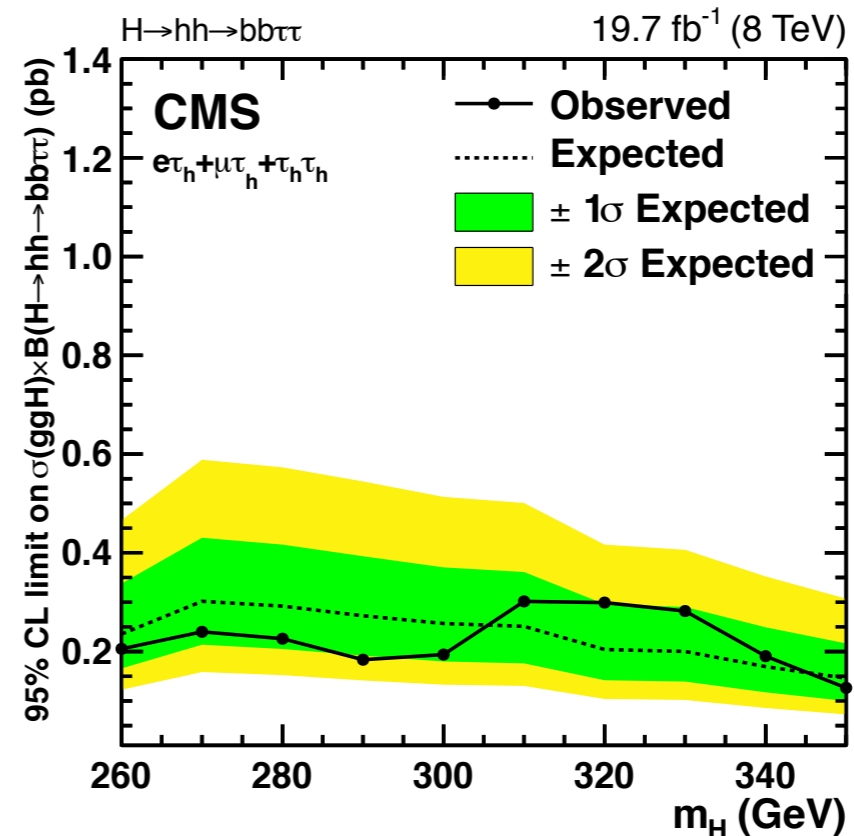
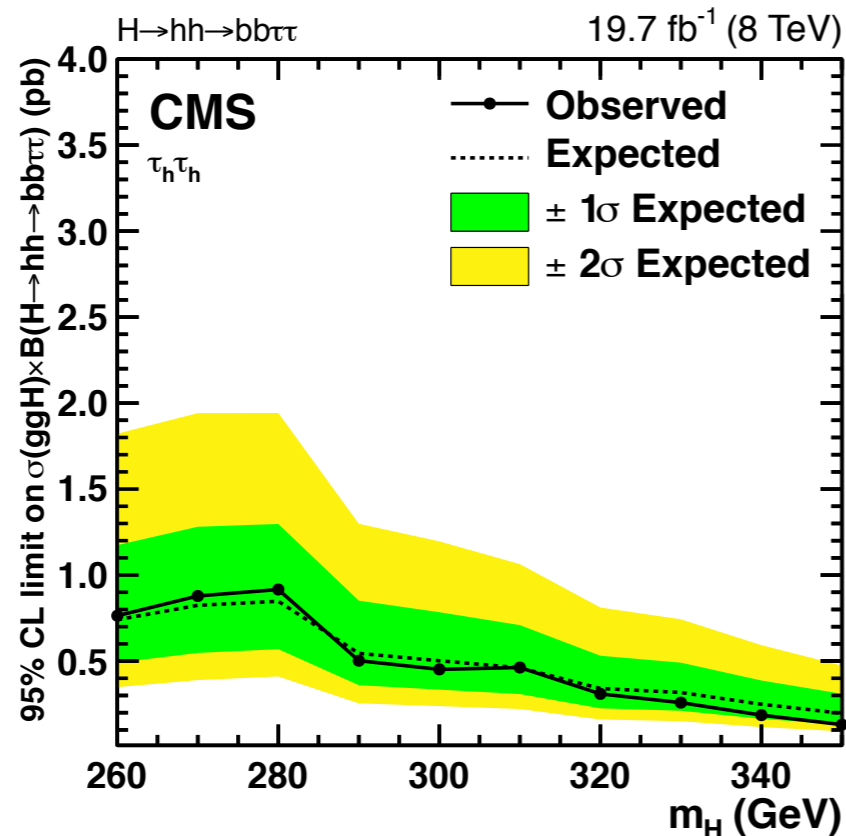
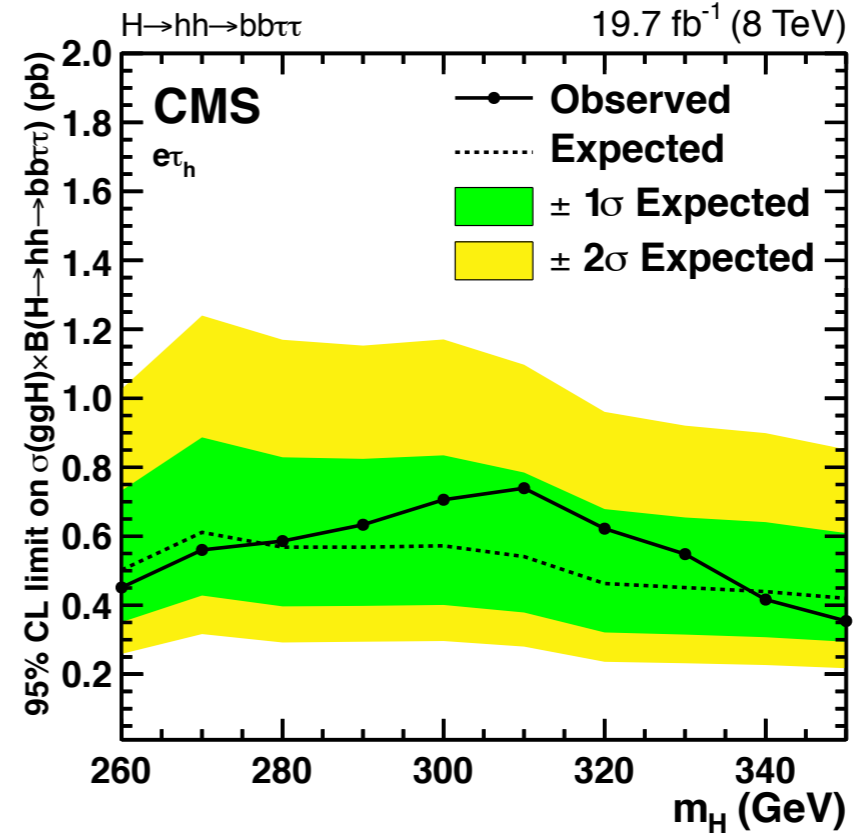
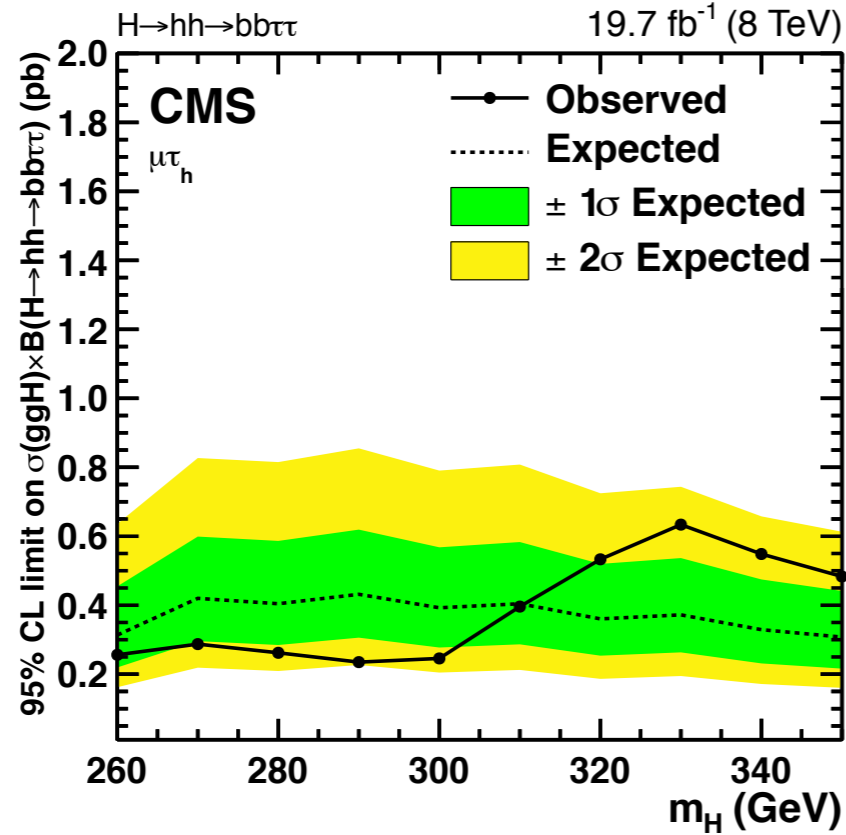


ATLAS: Only $\tau_\ell\tau_h$ channels
CMS: $\tau_\ell\tau_h$ and $\tau_h\tau_h$ channels.

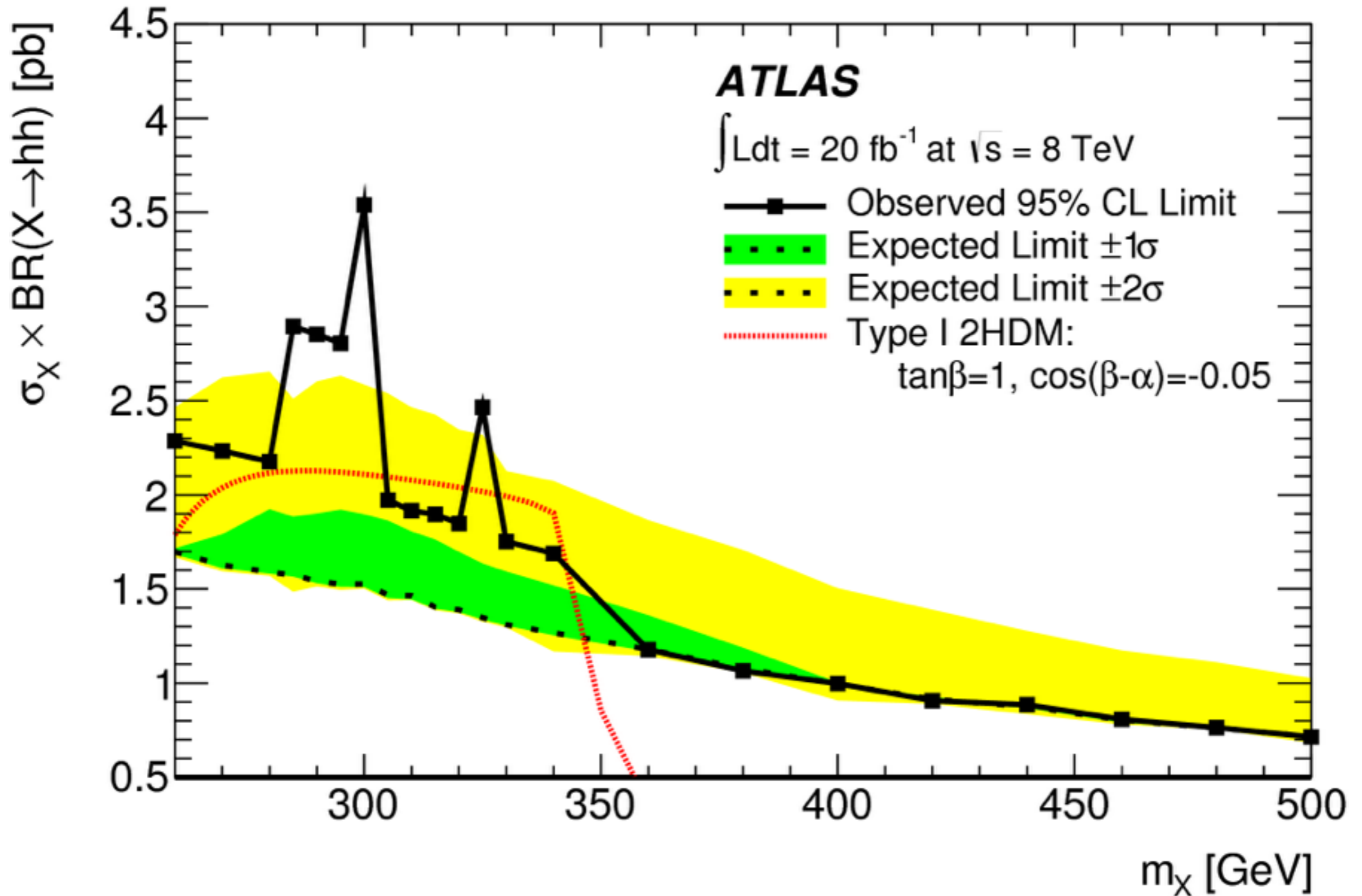
ATLAS: $\sigma \times \text{BR}(\text{H} \rightarrow \text{hh})$
CMS: $\sigma \times \text{BR}(\text{H} \rightarrow \text{hh} \rightarrow \text{bb}\tau\tau)$

Limit at 300 GeV:
 ATLAS: $3 \text{ pb} \times \text{BR}(\text{hh} \rightarrow \text{bb}\tau\tau) = 0.21 \text{ pb}$
 CMS: 0.2 pb

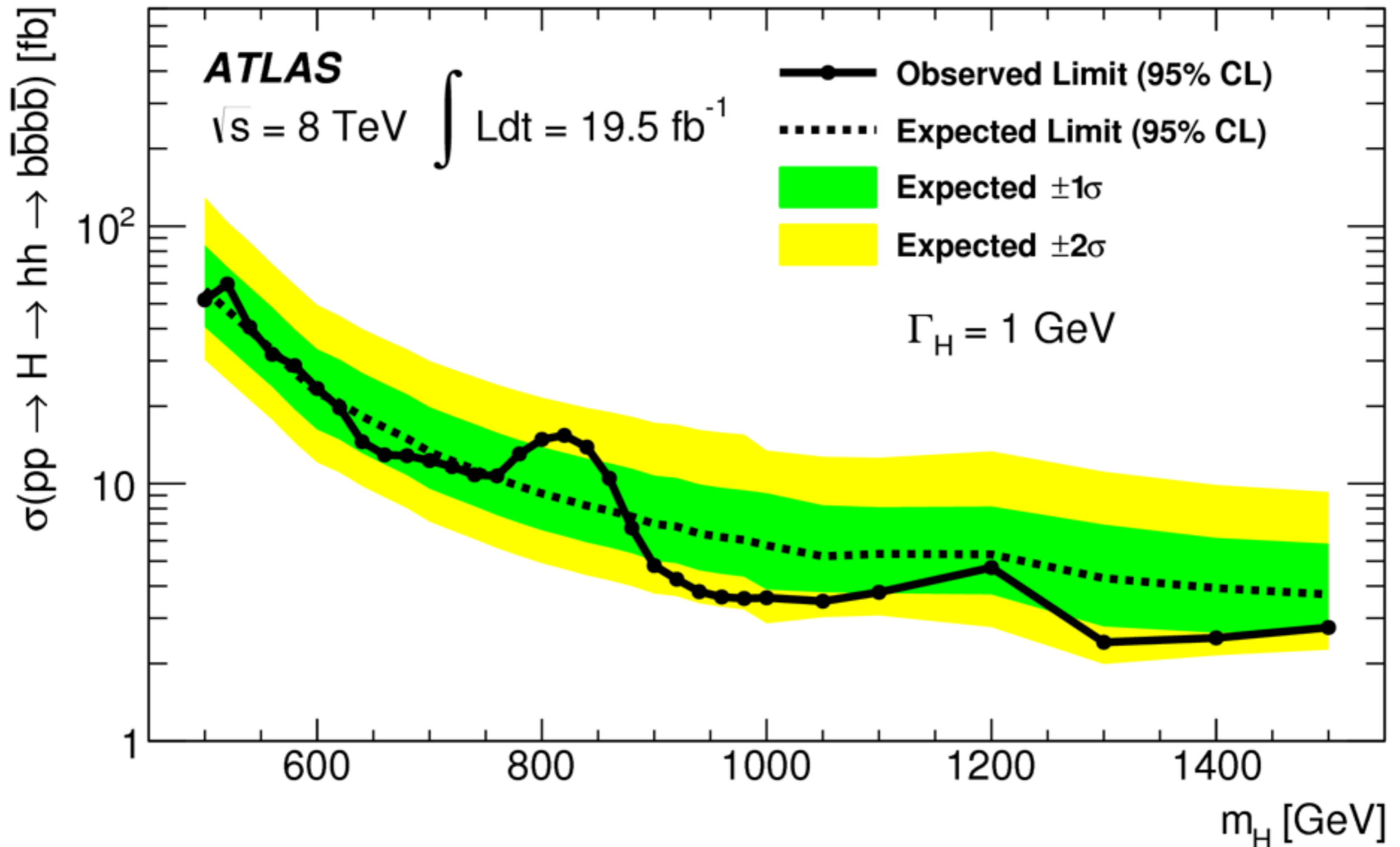
CMS $hh \rightarrow bb\tau\tau$ Results



ATLAS $H \rightarrow hh \rightarrow b\bar{b}\gamma\gamma$



ATLAS $H \rightarrow hh \rightarrow b\bar{b}b\bar{b}$



2HDM Models

- Higgs sector of 2HDM models described by parameters:
 - ▶ 4 Higgs masses
 - ▶ $\tan\beta$ (ratio of vacuum expectation values, vev)
 - ▶ α (mixing between the two neutral CP even states h, H).
- Several different ‘types’ of 2HDM:
 - ▶ Type I: One doublet couples to V (“fermiophobic”), one to fermions.
 - ▶ Type II: “MSSM like” model, one doublet couples to up-type quarks, one to down-type quarks.
 - ▶ Type III: “Lepton-specific” model, Higgs bosons have same couplings to quarks as type I and to leptons as in type II.
 - ▶ Type IV: “Flipped” model, Higgs bosons have same couplings to quarks as in type II and to leptons as in type I.
- For more specific MSSM models m_h fully determined at tree level by m_A and $\tan\beta$.

Prospects for Measuring λ_{HHH} at the HL-LHC

Signal Significance	had-had	lep-had (e channel)	lep-had (μ channel)	lep-lep (ee channel)	lep-lep ($\mu\mu$ channel)	lep-lep ($e\mu$ channel)
S/\sqrt{B} ($\lambda = 0\lambda_{\text{SM}}$)	0.66	0.49	0.45	0.044	0.055	0.091
S/B ($\lambda = 0\lambda_{\text{SM}}$)	0.023	0.023	0.023	0.00092	0.0016	0.0017
S/\sqrt{B} ($\lambda = 1\lambda_{\text{SM}}$)	0.47	0.35	0.33	0.033	0.036	0.062
S/B ($\lambda = 1\lambda_{\text{SM}}$)	0.016	0.016	0.017	0.00069	0.0010	0.0012
S/\sqrt{B} ($\lambda = 2\lambda_{\text{SM}}$)	0.29	0.25	0.23	0.023	0.025	0.044
S/B ($\lambda = 2\lambda_{\text{SM}}$)	0.010	0.012	0.011	0.00048	0.00074	0.00084
S/\sqrt{B} ($\lambda = 10\lambda_{\text{SM}}$)	1.0	0.56	0.52	0.062	0.072	0.14
S/B ($\lambda = 10\lambda_{\text{SM}}$)	0.036	0.027	0.026	0.0013	0.0021	0.0027

The Unbearable Lightness of M_H ...

Observed mass
(~125GeV)

$$M_H^2 = M_{\text{bare}}^2 + \left(\text{Higgs self-energy loop} \right) + \left(\text{top quark loop} \right) + \left(\text{W/Z loop} \right)$$

Bare mass to cancel radiative corrections

Radiative corrections, top loop dominates: $\sim m_t^2 \Lambda^2$
 Λ^2 : the energy scale at which the SM breaks down

Need a cancellation to 33 digits if Λ is at the Planck scale ($\sim 10^{19}$ GeV) - fine tuning!

Very strong motivation for new physics at TeV scale!